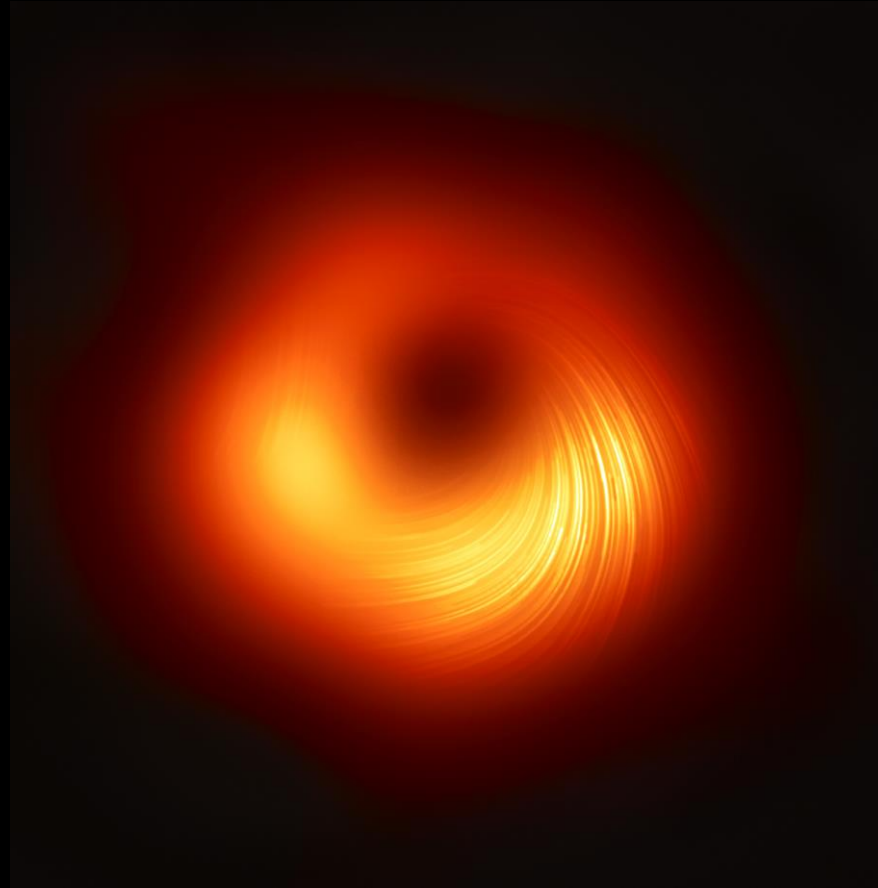


M87* in Polarized Light

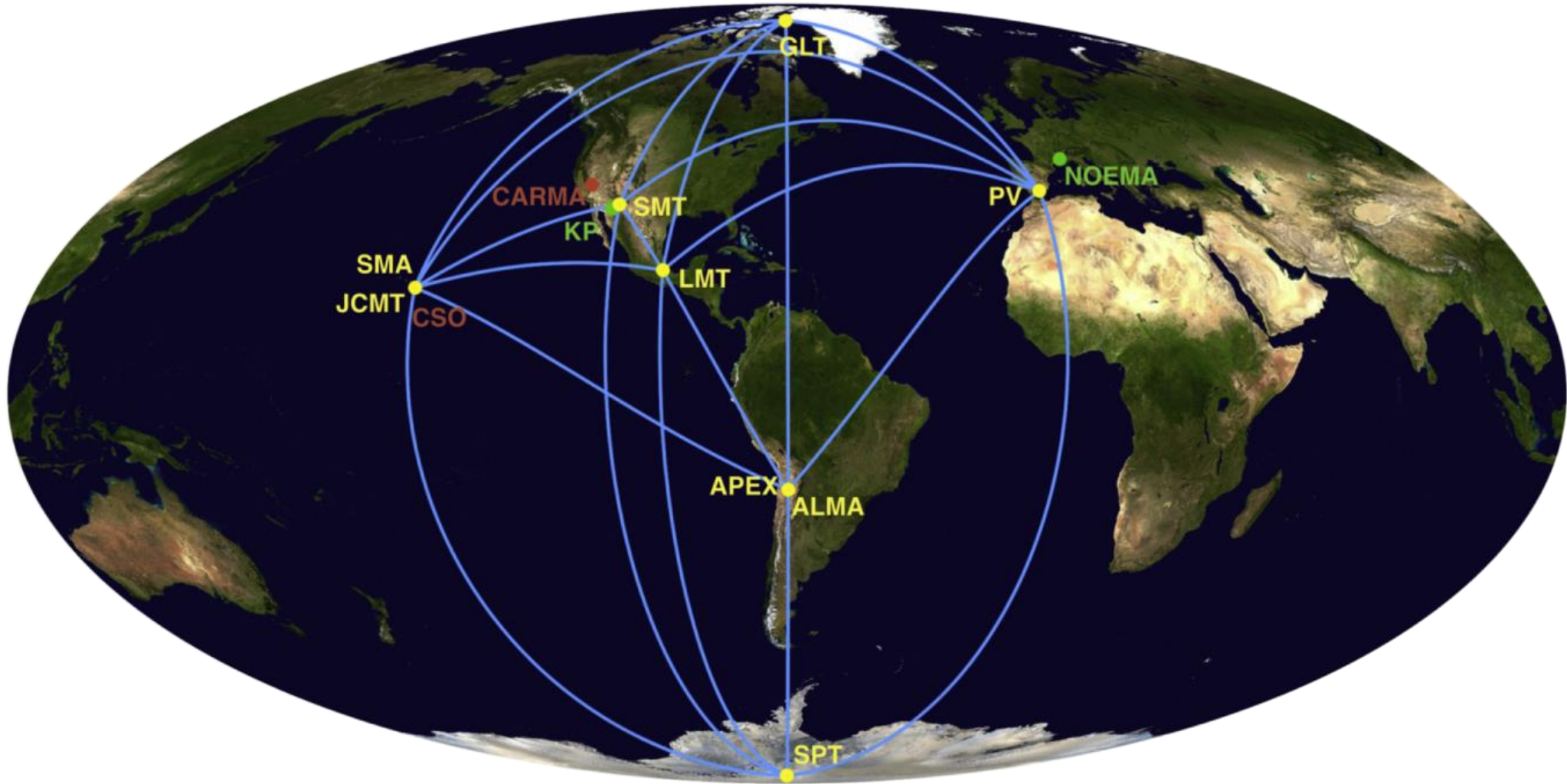
Andrew Chael

Princeton

November 15, 2023



The Event Horizon Telescope: Instrument



$$\text{Resolution} \approx \frac{\lambda}{d_{\text{Earth}}} \approx \frac{1.3 \text{ mm}}{1.3 \times 10^{10} \text{ mm}} \approx 20 \mu\text{as}$$

The Event Horizon Telescope: People



300+ members

60 institutes

20 countries

from Europe, Asia, Africa,
North and South America.

EHTC Paper VII,VIII,IX writing teams

Monika Mościbrodzka



Iván Martí-Vidal



Sara Issaoun



Jongho Park



Maciek Wielgus



Angelo Ricarte



Jason Dexter



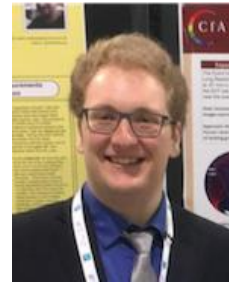
Andrew Chael



Alejandra Jiménez-Rosales



Daniel Palumbo



Dom Pesce



John Wardle



Svetlana Jorstad



Ioannis Myserlis



Freek Roelofs



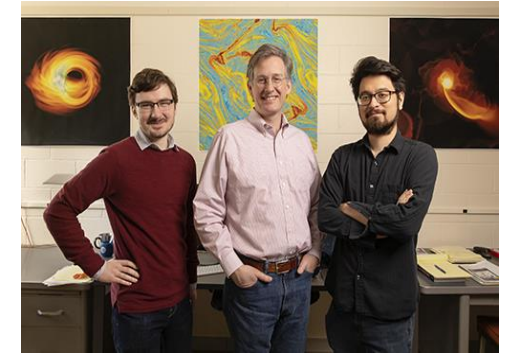
Abhishek Joshi



Avery Broderick



**Ben Prather, Charles Gammie,
George Wong**



M87*

$$M_{BH} = (6.5 \pm 0.7) \times 10^9 M_{\odot}$$

$$D = (16.8 \pm 0.8) \text{Mpc}$$

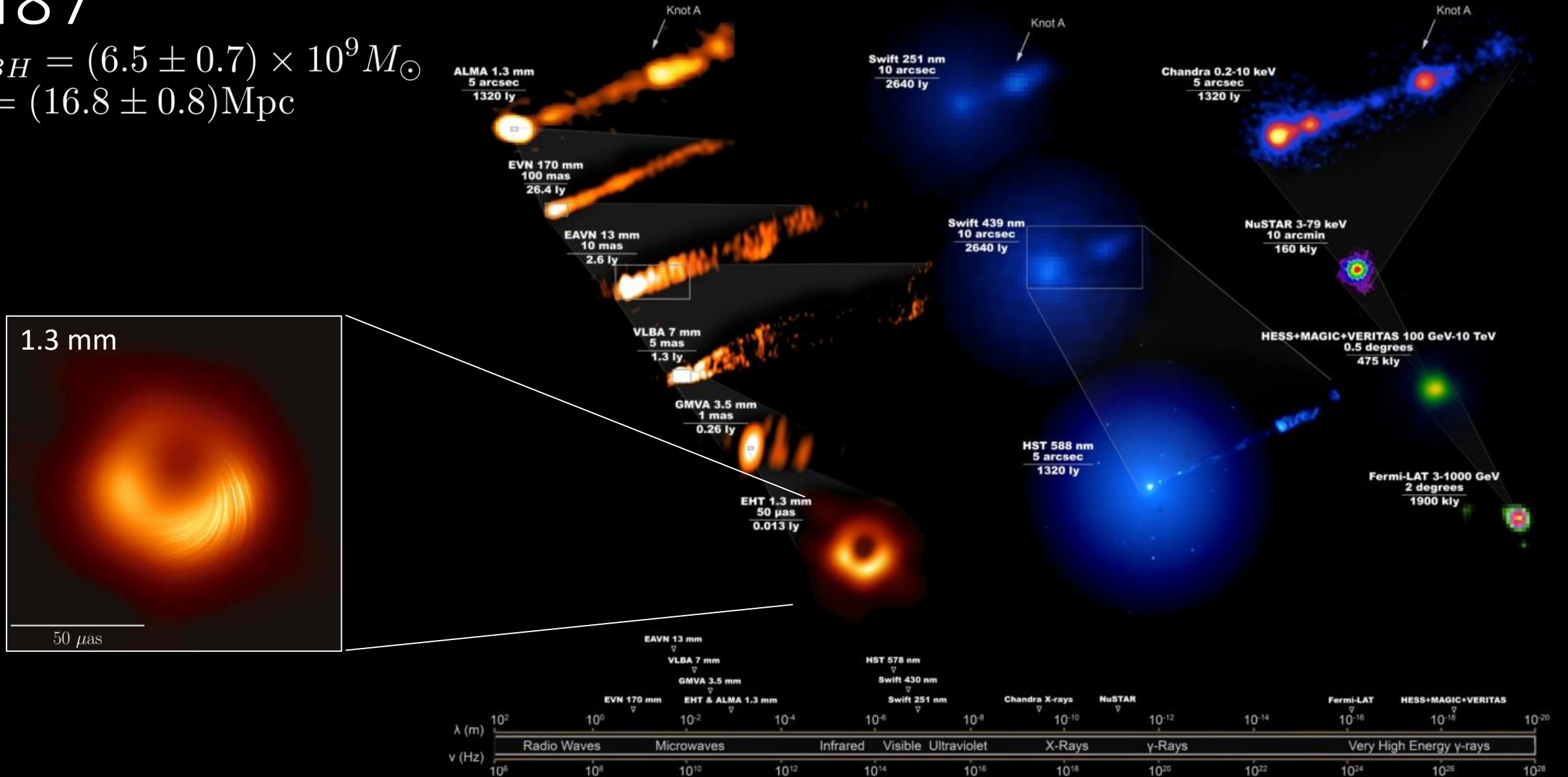
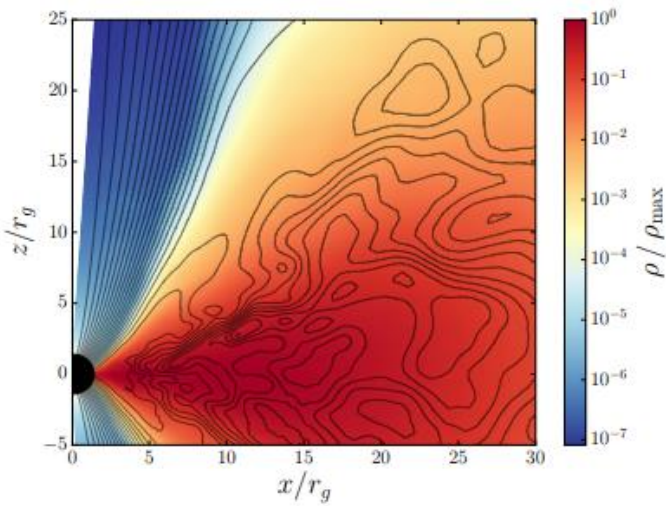


Image Credit: The EHT Multi-wavelength Science Working Group; the EHT Collaboration; ALMA (ESO/NAOJ/NRAO); the EVN; the EAVN Collaboration; VLBA (NRAO); the GMVA; the Hubble Space Telescope; the Neil Gehrels Swift Observatory; the Chandra X-ray Observatory; the Nuclear Spectroscopic Telescope Array; the Fermi-LAT Collaboration; the H.E.S.S. collaboration; the MAGIC collaboration; the VERITAS collaboration; NASA and ESA. Composition by J. C. Algaba

What is the magnetic field structure close to the horizon?

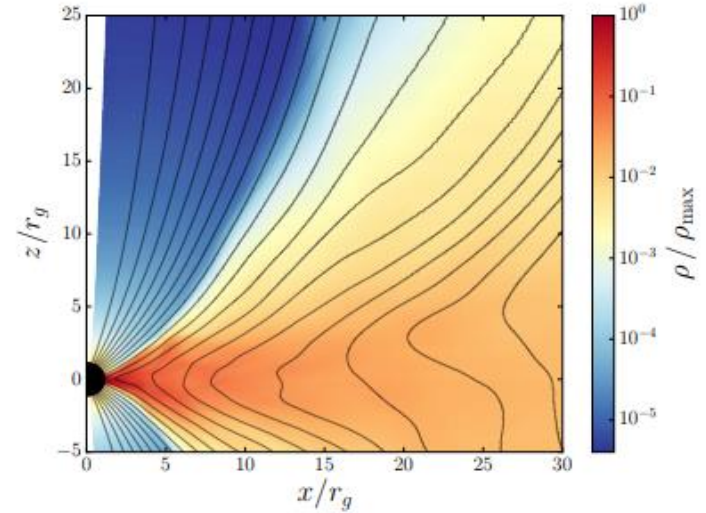
Two accretion states that depend on the accumulated magnetic flux on horizon

Magnetic fields are weak and turbulent



“SANE”

Strong, coherent magnetic fields build up on the horizon



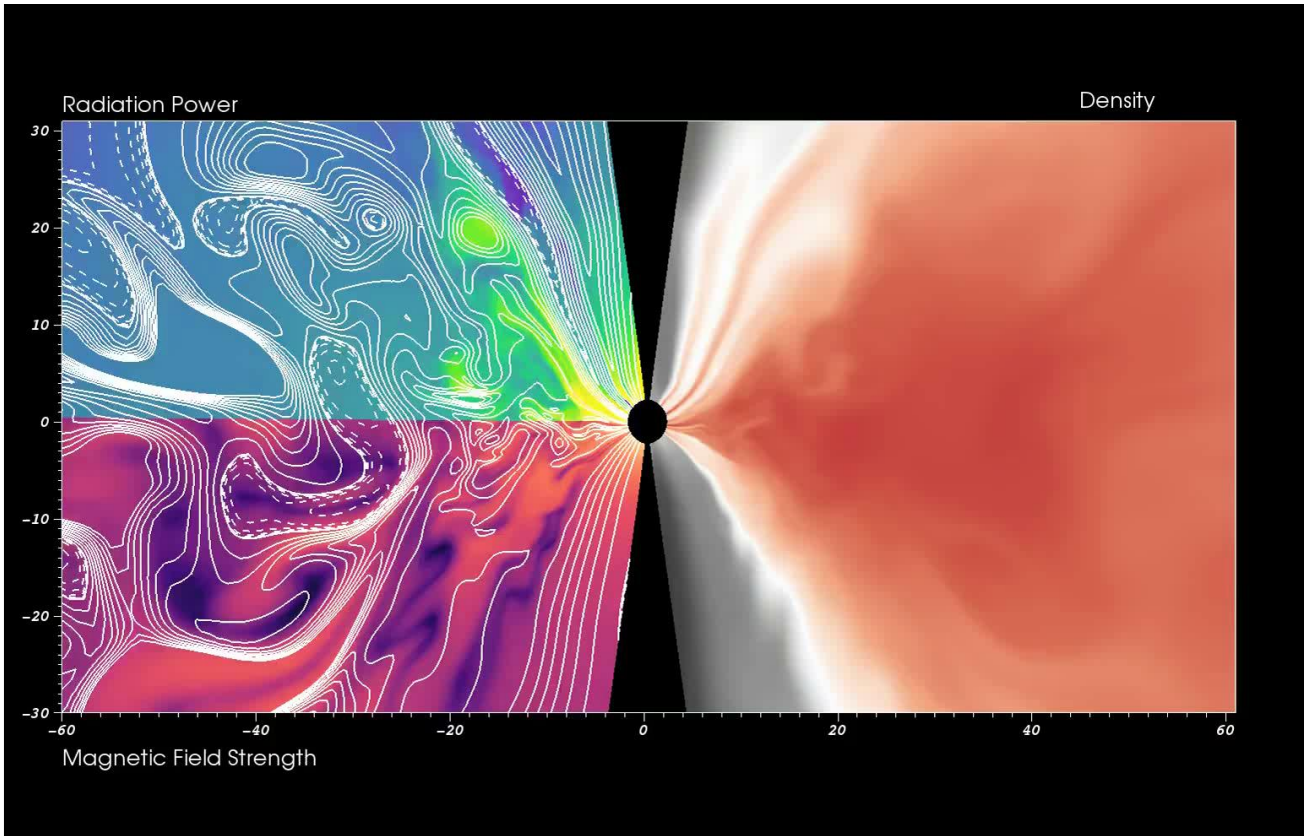
“MAD” - Magnetically Arrested Disk

Note: ‘strong’ fields mean dynamically important ones → ~10 G at the horizon for M87

Blandford-Znajek (1977): $P_{\text{jet}} \propto \Phi_B^2 a^2$

↑ magnetic flux ↘ BH spin

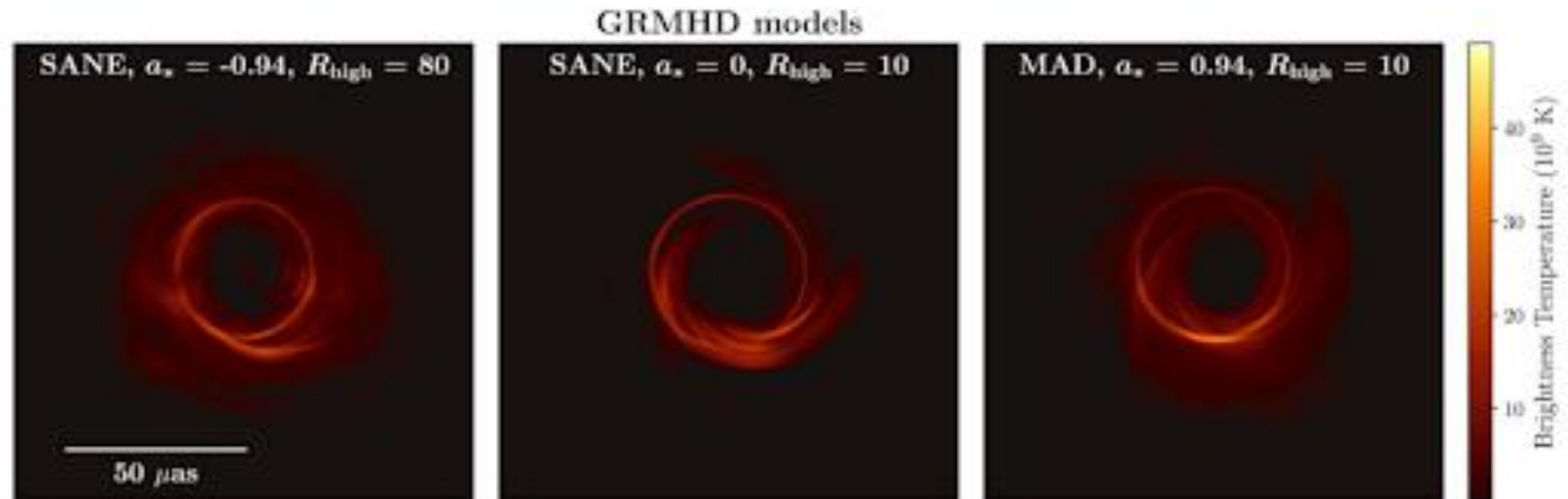
Interpreting Images with GRMHD Simulations



- GRMHD simulations of **radiatively inefficient disks** are the primary theoretical tool for interpreting EHT images.
- Hot ($10^{10} < T < 10^{12}$ K), dilute ($10^4 < n < 10^7$ cm⁻³), magnetized ($1 \text{ G} < |B| < 50 \text{ G}$) plasma naturally satisfies constraints on
 - Image brightness
 - Faraday Rotation / low linear polarization
 - Faraday Conversion / low circular polarization
- GRMHD simulations naturally **couple the accretion disk, black hole, and jet**
 - Jet launching in simulations is universal and driven by BH spin

Scoring GRMHD Simulations: before polarization

- **Most simulation models can be made to fit total intensity observations alone by tweaking free parameters (mass, PA, total flux density)**



- An additional constraint on **jet power** ($\geq 10^{42}$ erg/sec) rejects all spin 0 models
- Can we do better with polarization?

Outline

1. How do we obtain a polarized image of M87* with the EHT?
2. How do we interpret the polarized image of M87*?
3. Connection between polarized images and EM energy flux



First M87 Event Horizon Telescope Results. VII. Polarization of the Ring

The Event Horizon Telescope Collaboration
(See the end matter for the full list of authors.)

Received 2020 November 23; revised 2021 February 15; accepted 2021 February 16; published 2021 March 24

First M87 Event Horizon Telescope Results. VIII. Magnetic Field Structure near The Event Horizon

The Event Horizon Telescope Collaboration
(See the end matter for the full list of authors.)

Received 2020 December 2; revised 2021 February 3; accepted 2021 February 8; published 2021 March 24

First M87 Event Horizon Telescope Results. IX. Detection of Near-horizon Circular Polarization

The Event Horizon Telescope Collaboration
(See the end matter for the full list of authors.)

Received 2023 June 19; revised 2023 September 5; accepted 2023 September 12; published 2023 November 8

Black Hole Polarimetry I. A Signature of Electromagnetic Energy Extraction

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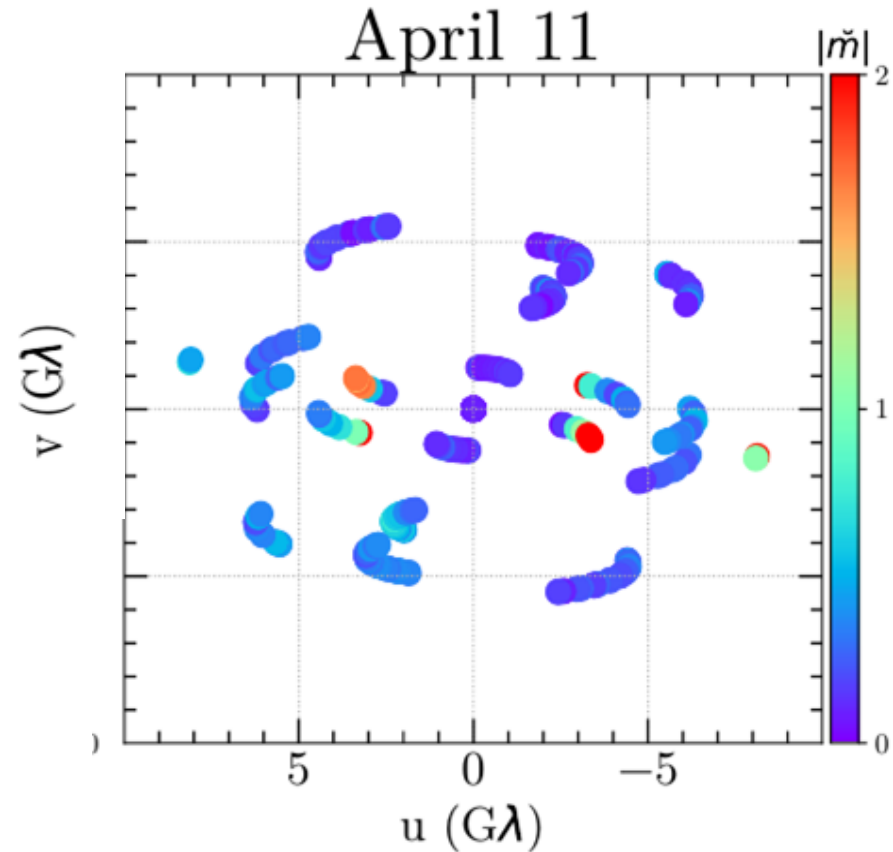
⁴Department of Astrophysical Sciences, Princeton University, Princeton, NJ 08544, USA

Received 2023 July 12; revised 2023 August 11; accepted 2023 September 11; published 2023 November 14

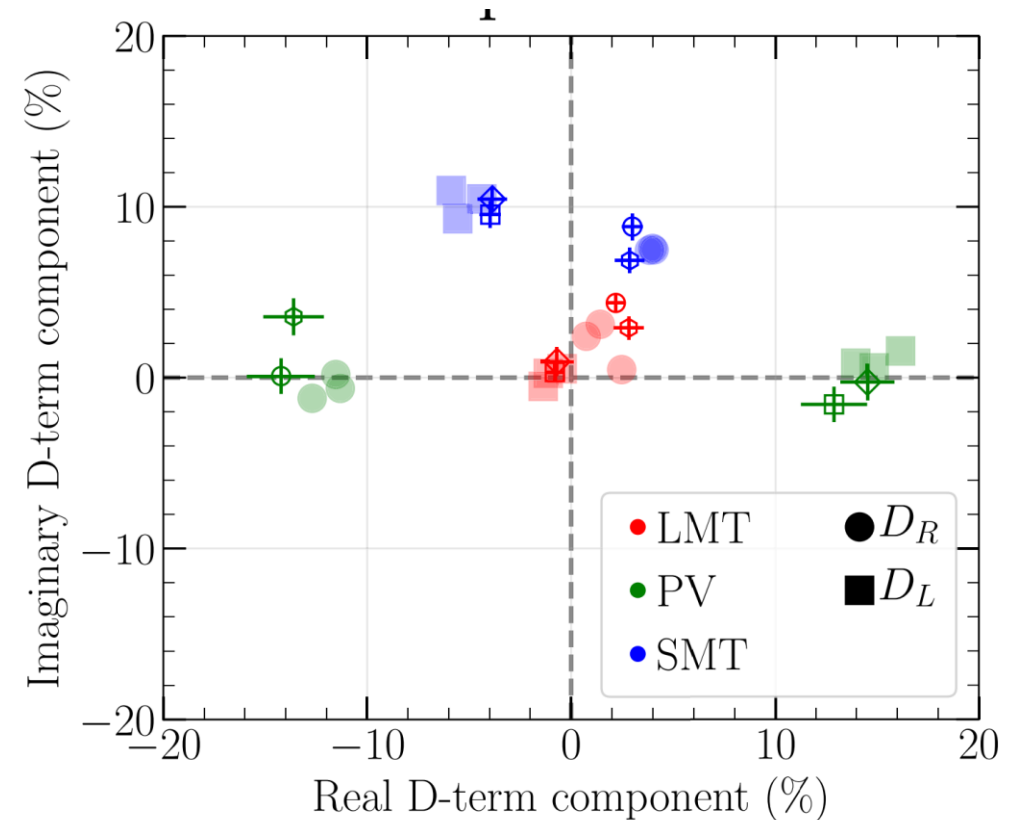
How do we obtain a polarized image of M87* with the EHT?

Two Challenges of EHT polarimetric imaging

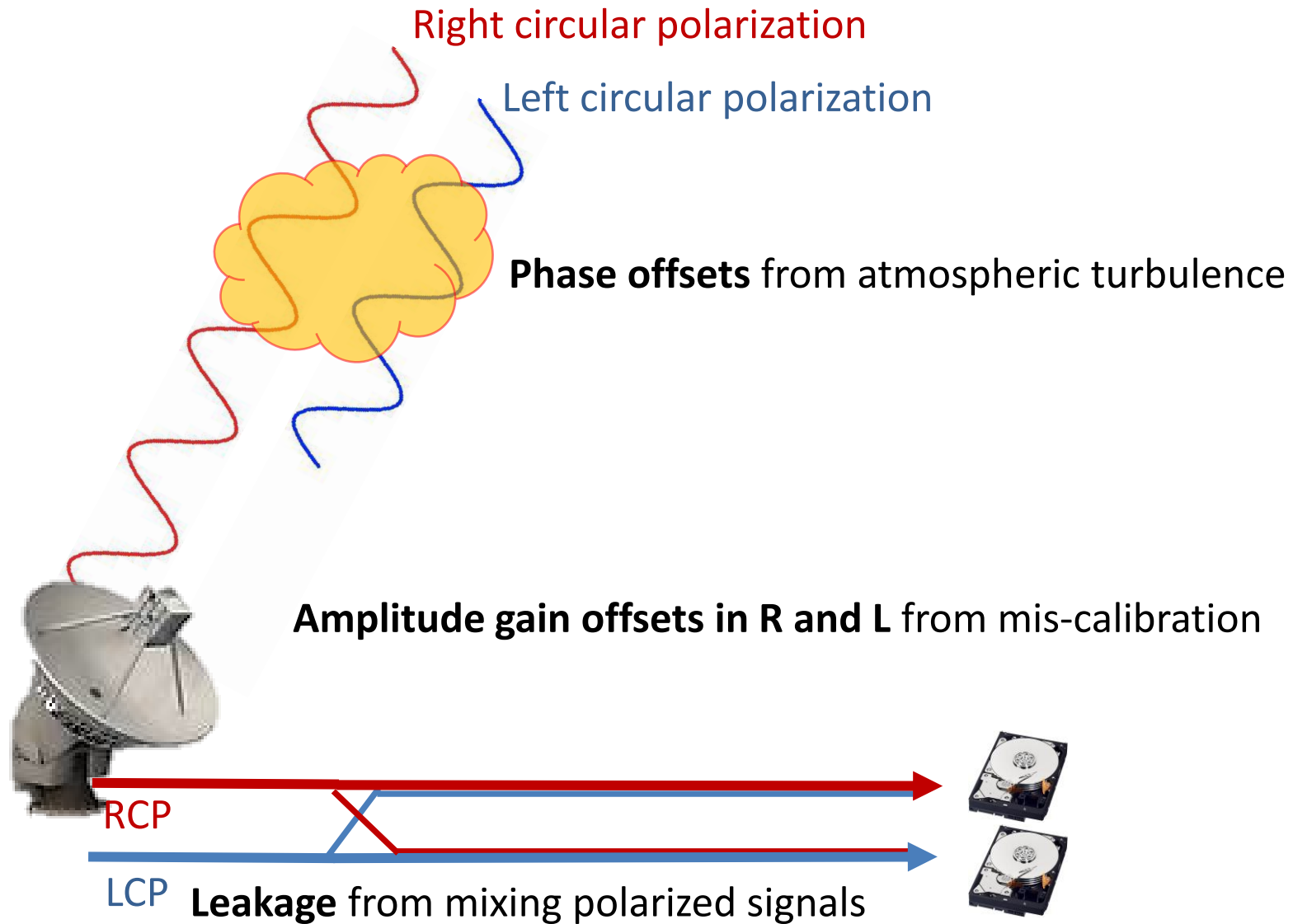
1. EHT coverage is **sparse**: inversion of image from the data is highly unconstrained



Data at each station are corrupted by unknown polarimetric **leakage** and polarization-dependent complex gain factors



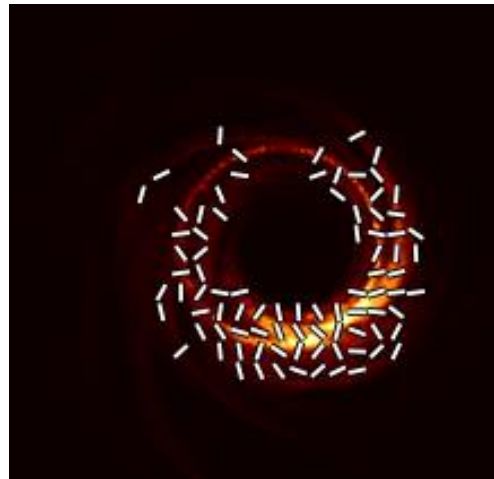
Corrupting effects at EHT stations



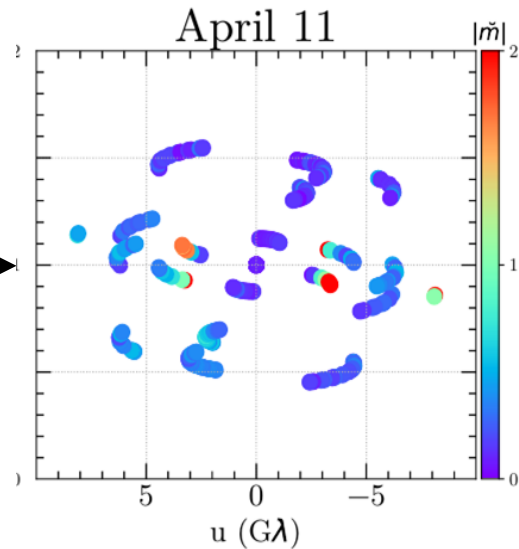
Data at each station are corrupted by unknown polarimetric **leakage** and **complex gain** factors

Solving for the Image

True Image

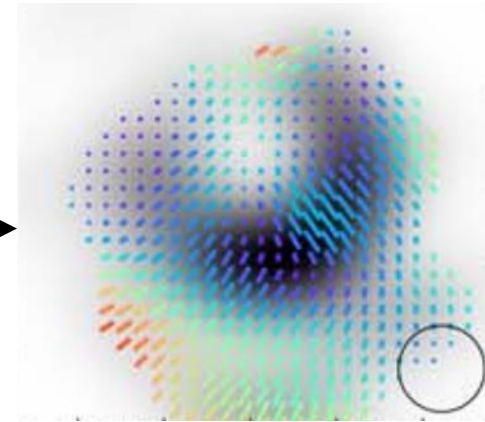


Sparse/ Corrupted
Measurements



RECONSTRUCTION
ALGORITHM

Reconstruction



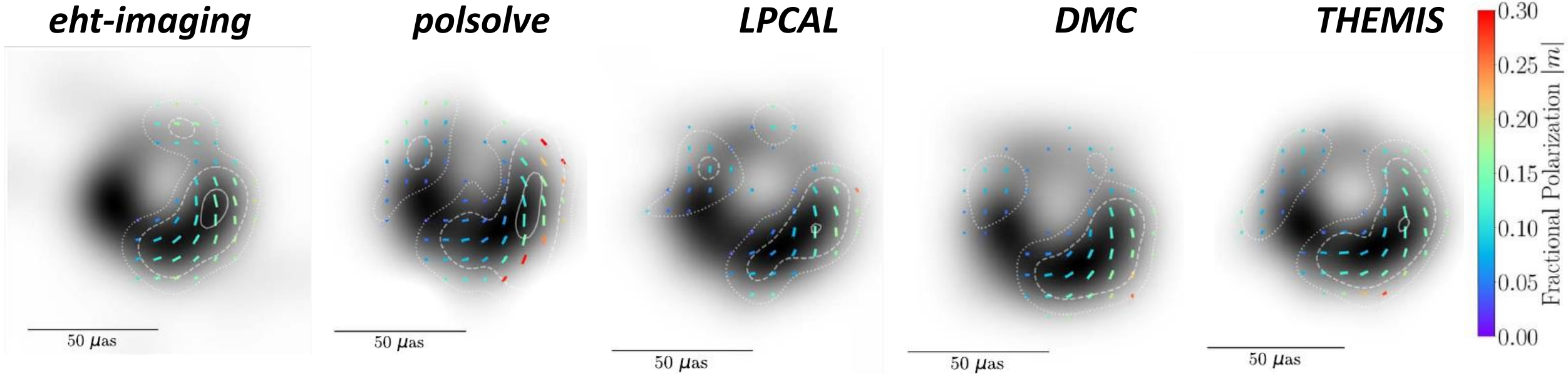
Several different types of reconstruction algorithms now used:

- **CLEAN-based**: standard and efficient, but can have difficulties on very sparse data
 - LPCAL/GPCAL (Park+ 2021) and polysolve (Marti-Vidal+ 21)
- **Regularized Maximum Likelihood w/ Gradient Descent**: fast and flexible, but lots of hyperparameters
 - eht-imaging (Chael+ 2016, 2018)
- **Bayesian MCMC posterior exploration**: fully characterizes uncertainty, but expensive
 - Themis (Broderick+ 21), DMC (Pesce+ 21)

credit: Katie Bouman, Andrew Chael,

EHTC 2021. Paper VII

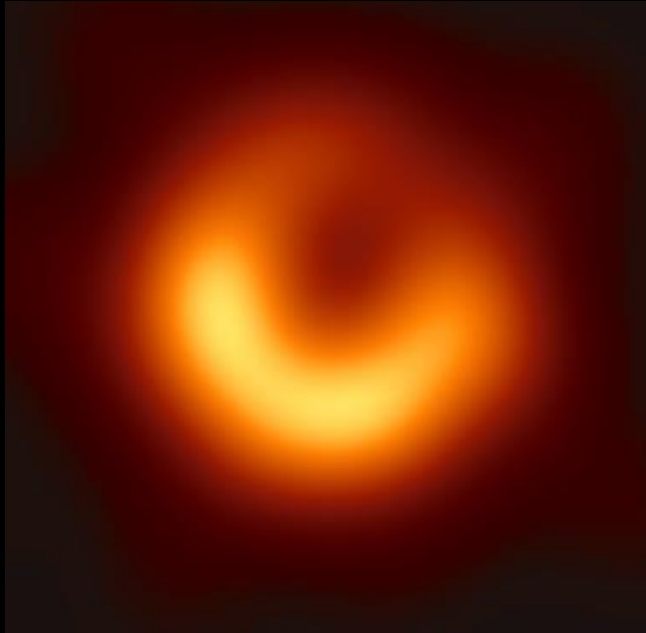
Linear Polarization Images from five vetted methods



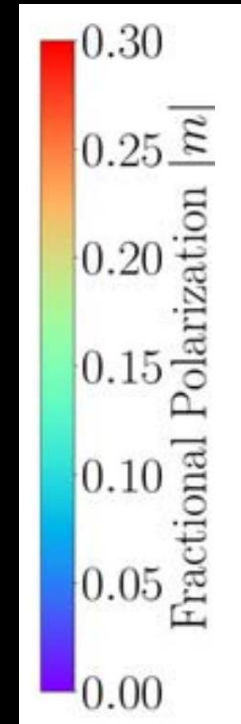
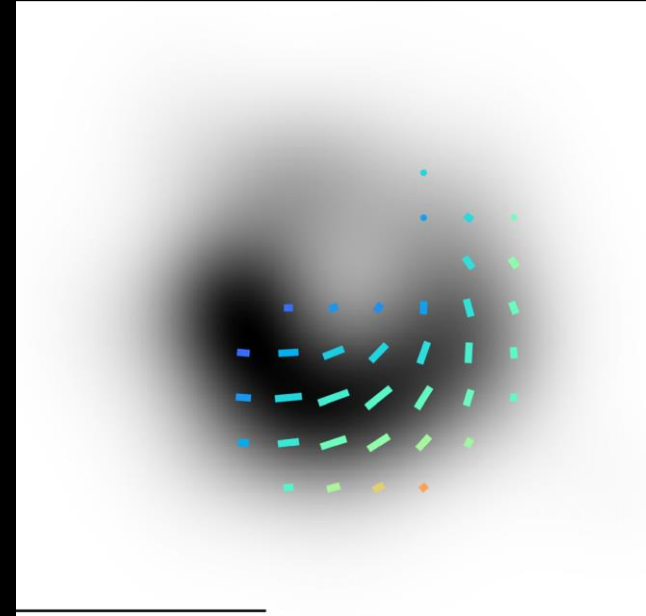
- All methods show similar polarization structure
- Polarization is concentrated in the southwest
- Polarization angle structure is predominantly **azimuthal**
- Overall level of polarization is **somewhat weak**, $|m|$ rises to $\sim 15\%$

M87* in linear polarization

Total intensity



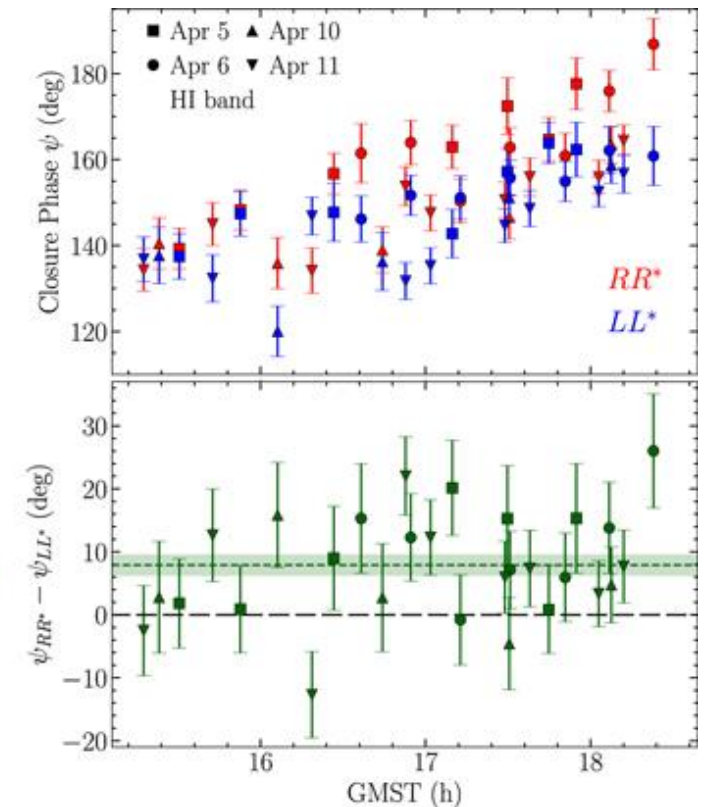
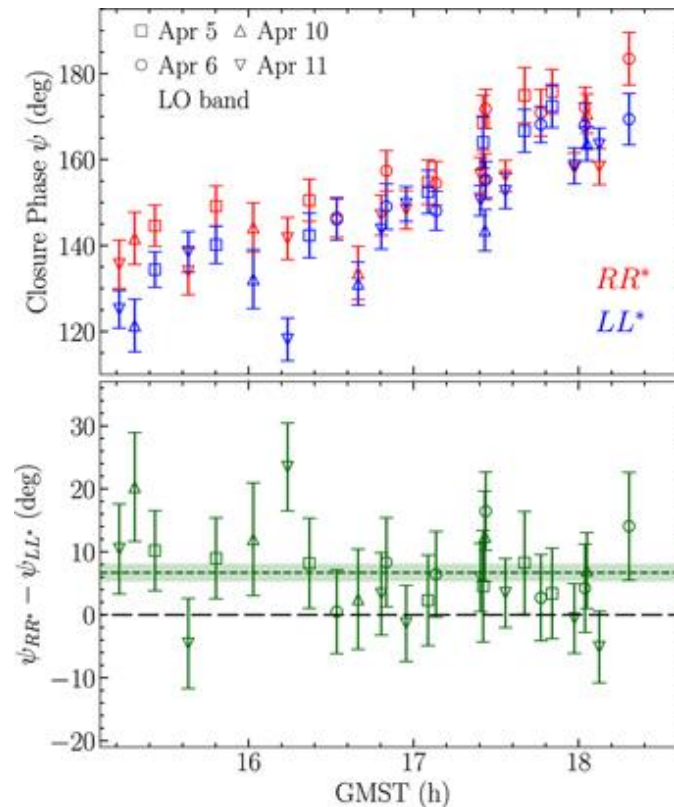
Linear Polarization



- Polarization is concentrated in the southwest
- Polarization angle structure is predominantly **helical**
- Overall level of polarization is **somewhat weak**, ~15 %

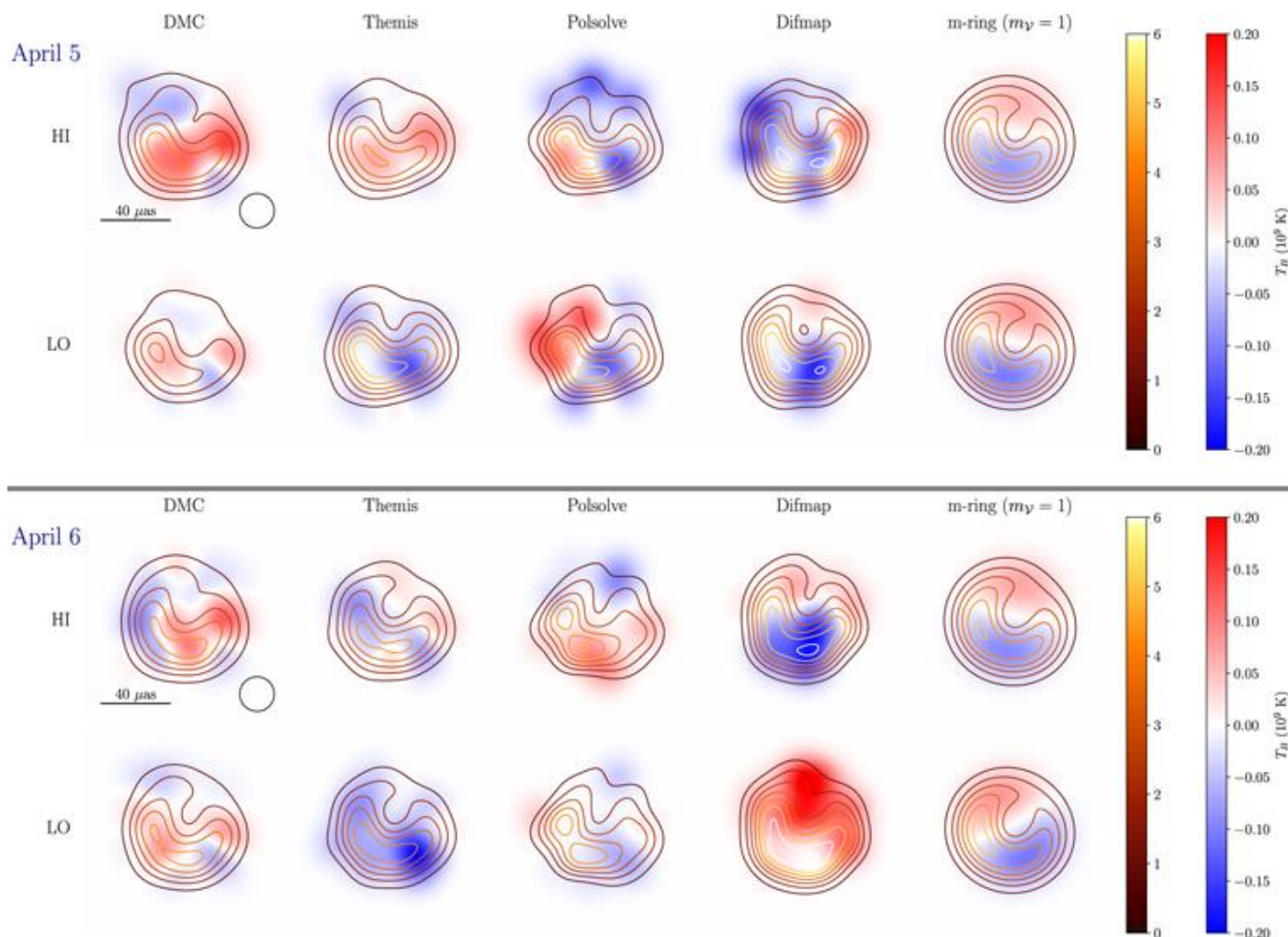
Horizon-Scale circular polarization is unambiguously detected by the EHT

- We detect an **offset** between closure phases in the RR and LL polarizations ($V=0.5(RR-LL)$)
- This is immune to relative gain offsets G_R / G_L
- Not seen on all triangle; upper limit of detected circular polarization in Fourier space is only 1%-10% of total intensity
- Can we constrain the image structure in circular polarization?



Horizon-Scale circular polarization *images* are **not** robustly recovered

- Different reconstruction methods make different assumptions about how to calibrate gains, D-terms, other systematics
- Methods do not show consistent Stokes V images
 - Not consistent between days
 - Not consistent between frequency bands
- Methods show a similar overall level of $|V|$ across the image
 - Use to place an upper limit on $\langle |v| \rangle < 3.7\%$

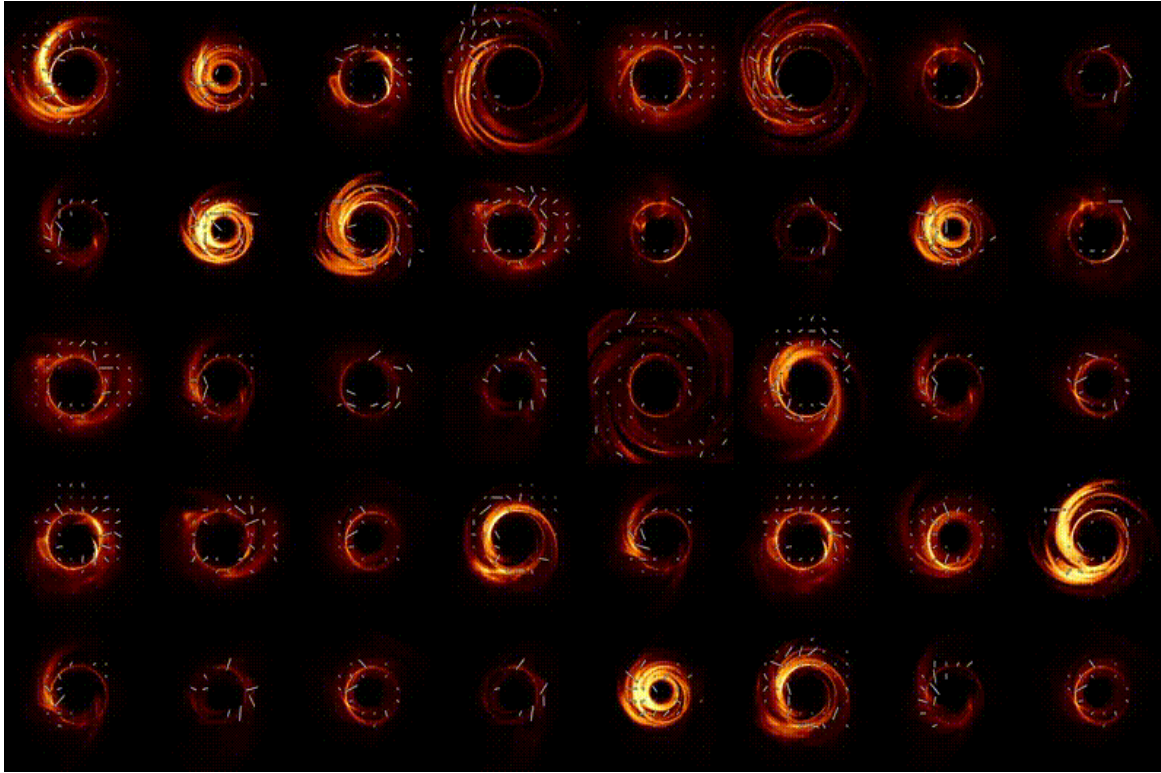


$$\langle |v| \rangle = \frac{\int |V/I| I dA}{\int I dA} \quad \text{Average resolved Circular Fraction}$$

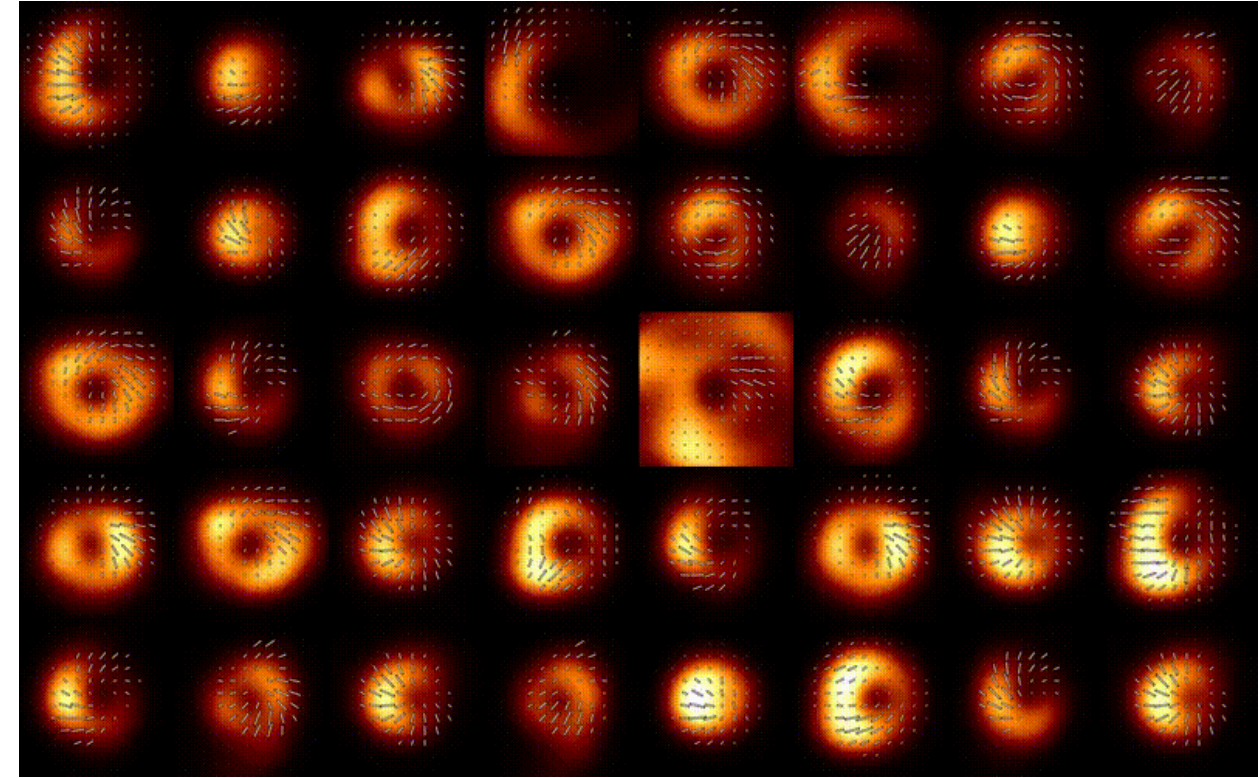
What do the EHT's polarization results tell us about the accretion flow?

GRMHD Simulation library

2 field states, 5 spins, >180k images



native resolution



EHT resolution

Images modeled with the ipole GRRT code (Moscibrodzka & Gammie 2018)

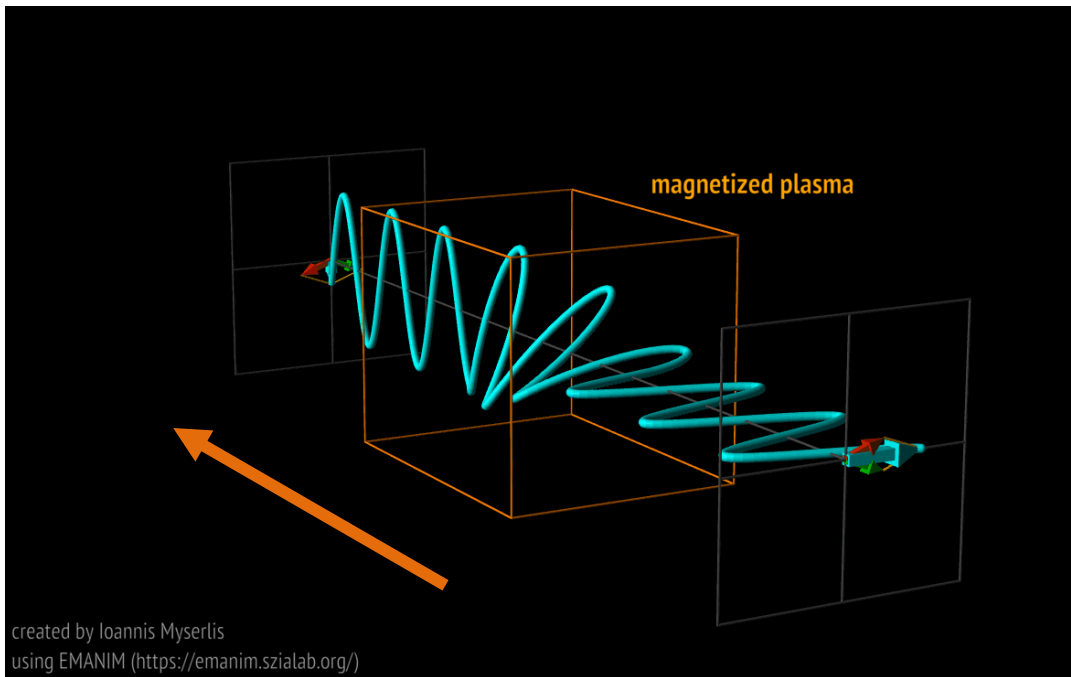
Two-temperature plasma model from Moscibrodzka et al. 2016

$$\frac{T_i}{T_e} = R_{\text{high}} \frac{\beta^2}{1 + \beta^2} + R_{\text{low}} \frac{1}{1 + \beta^2}$$

Two parameters set the electron temperature

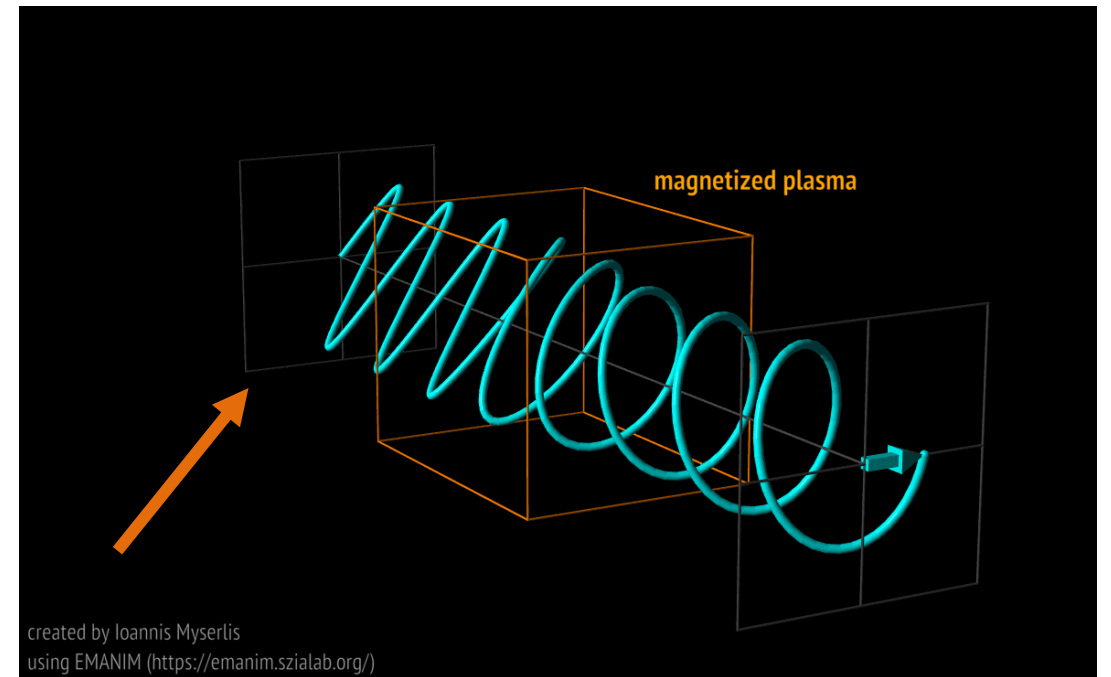
Faraday rotation and conversion are critical

Rotation



Field parallel to propagation matters

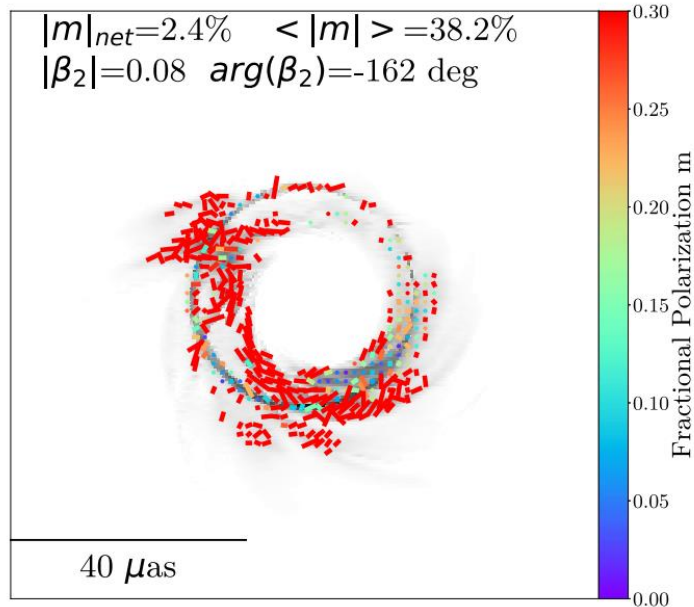
Conversion



Field parallel to linear polarization vector matters

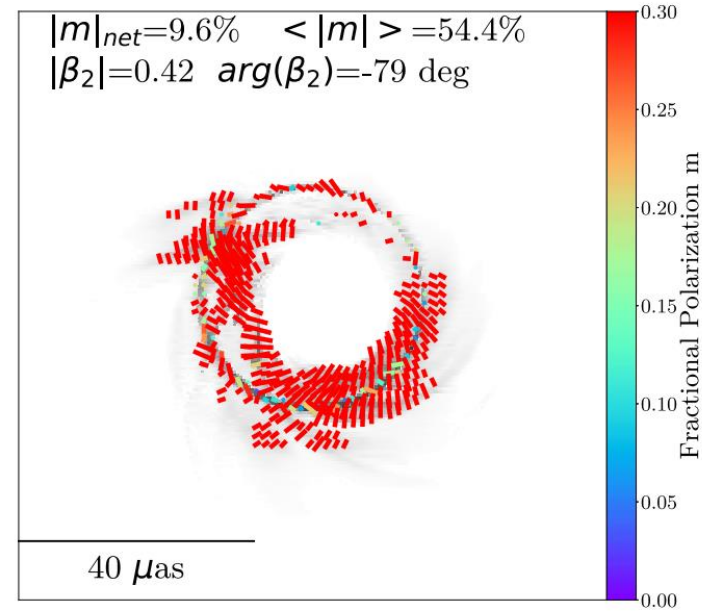
(Internal) Faraday rotation matters!

With rotation



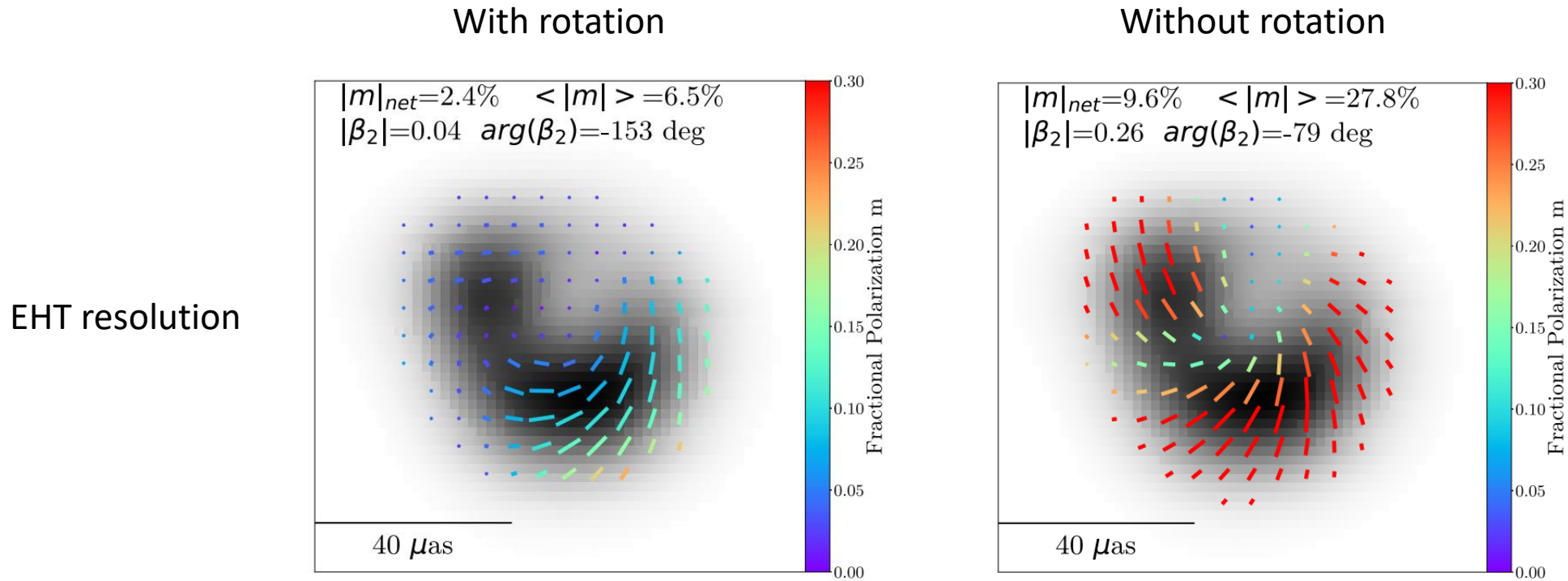
'infinite' resolution

Without rotation



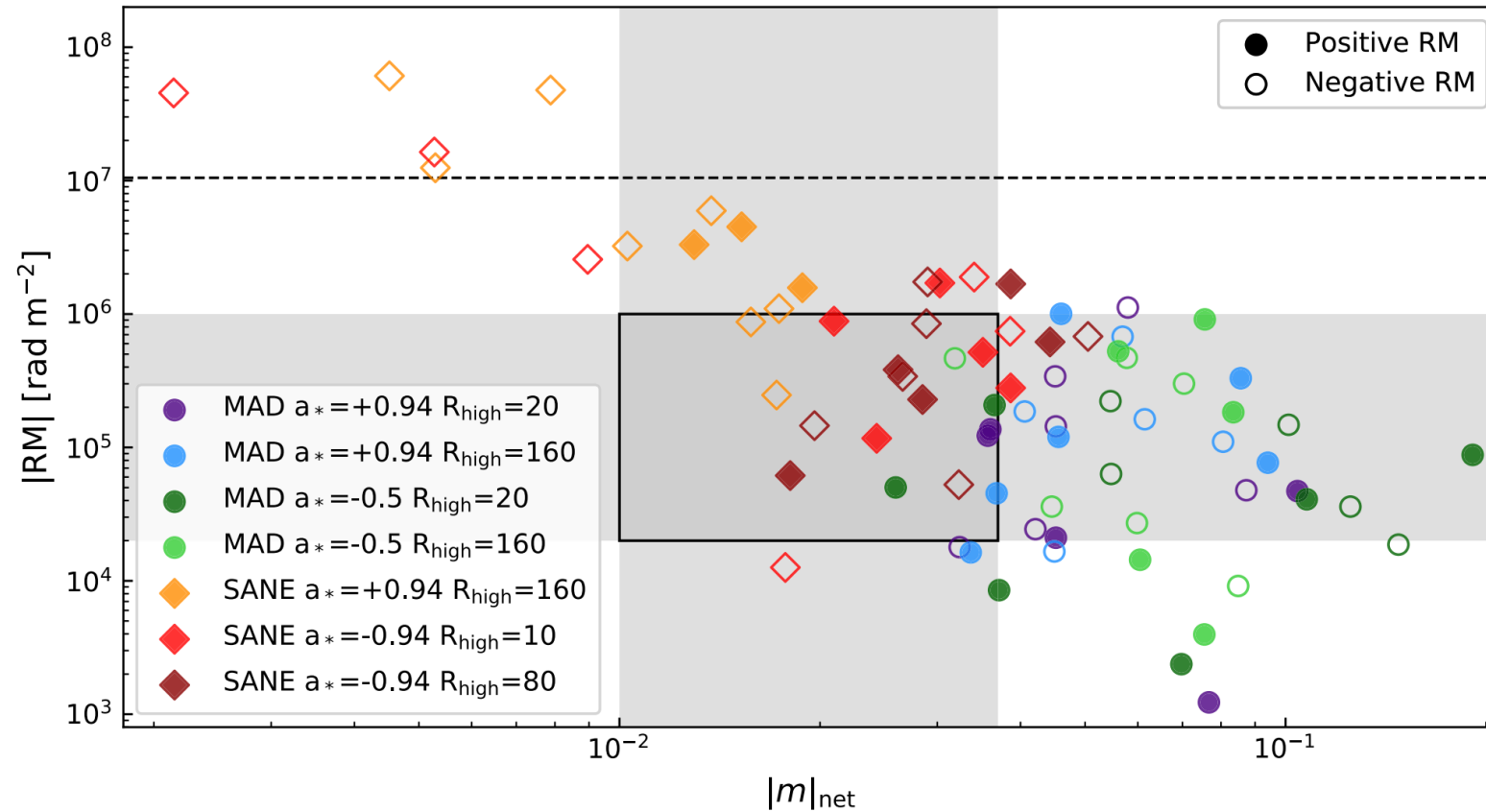
- Significant Faraday rotation on small scales
 - **scrambles** polarization directions
 - **depolarization** of the image when blurred to EHT resolution
 - **overall rotation** of the pattern when blurred to EHT resolution

(Internal) Faraday rotation matters!



- Significant Faraday rotation on small scales
 - **scrambles** polarization directions
 - **depolarization** of the image when blurred to EHT resolution
 - **overall rotation** of the pattern when blurred to EHT resolution

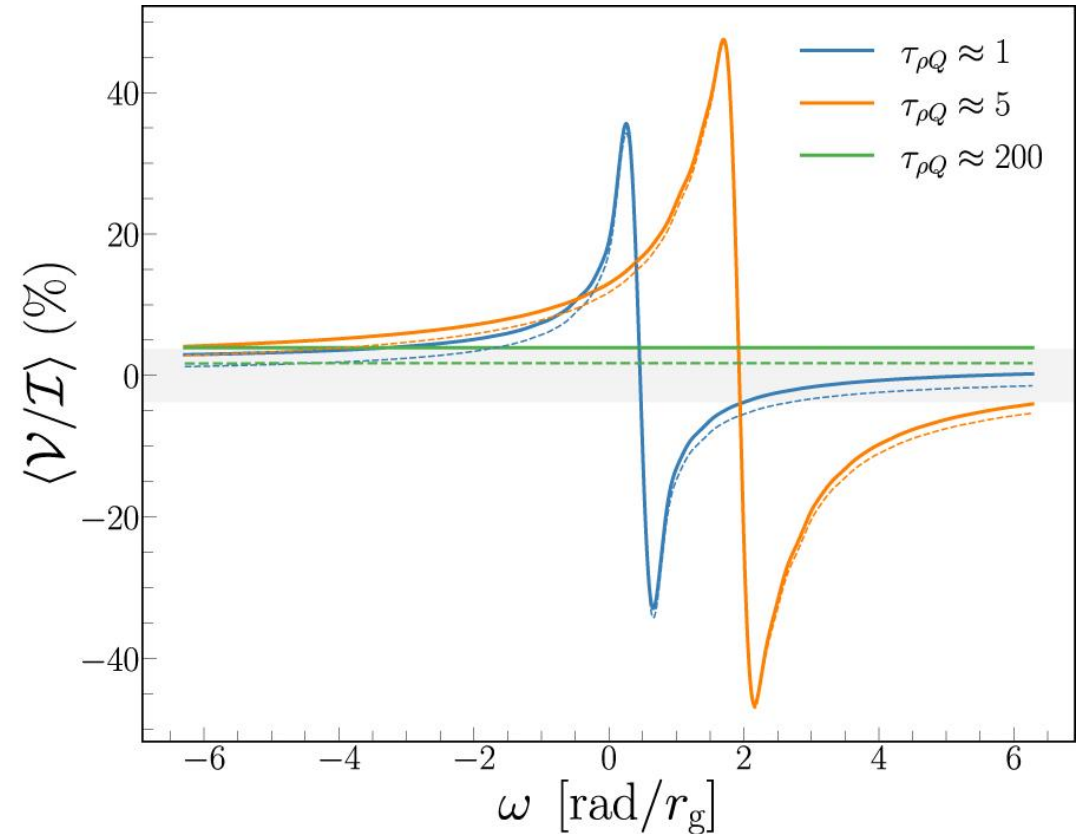
GRMHD simulations can explain M87's Rotation Measure



Important in future work to use simultaneous observations on larger scales to better constrain contributions of internal and any external Faraday rotation.

Most circular polarization is produced by **conversion**

- One-zone models and GRMHD simulations both confirm conversion is the dominant source of circular polarization in favored models
- In a uniform field geometry, Faraday conversion will typically produce more circular than linear polarization
- The interplay of conversion, rotation, and changing magnetic field direction along the line of sight determines the level and sign of circular polarization



Scoring simulations with polarization: Image metrics

Unresolved linear
polarization fraction

$$|m|_{\text{net}} = \frac{\sqrt{(\sum_i Q_i)^2 + (\sum_i U_i)^2}}{\sum_i I_i}$$

Unresolved circular
polarization fraction
(from ALMA)

$$|v|_{\text{net}} = \frac{|\sum_i V_i|}{\sum_i I_i}$$

Average resolved
linear fraction

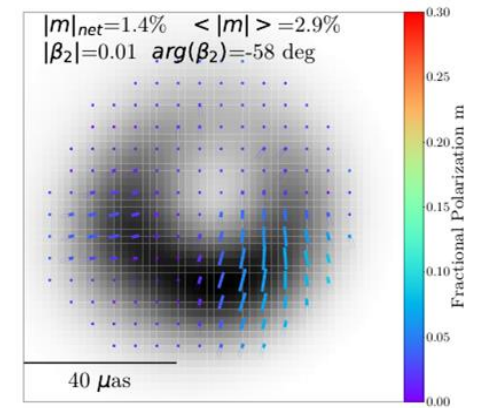
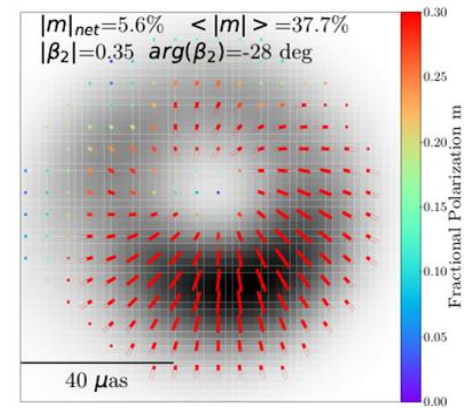
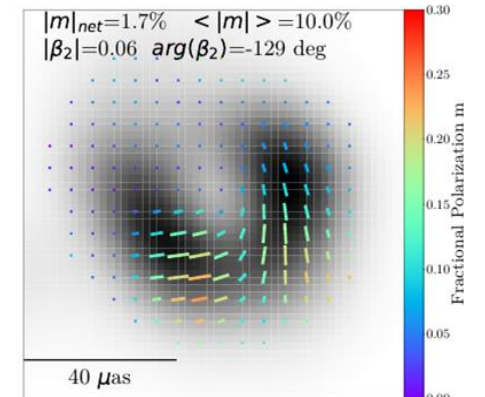
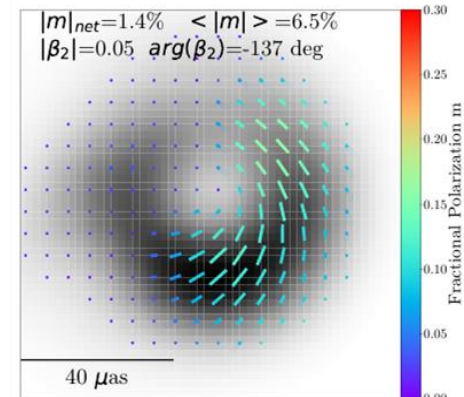
$$\langle |m| \rangle = \frac{\sum_i \sqrt{Q_i^2 + U_i^2}}{\sum_i I_i}$$

Average resolved
circular fraction

$$\langle |v| \rangle = \frac{\sum_i |V_i/I_i|}{\sum_i I_i}$$

Azimuthal Linear structure
2nd mode (Palumbo+ 2020)

$$\beta_2 = \frac{1}{I_{\text{ring}}} \int_{\rho_{\text{min}}}^{\rho_{\text{max}}} \int_0^{2\pi} P(\rho, \varphi) e^{-2i\varphi} \rho d\varphi d\rho$$



GRMHD images can be **strongly** or **weakly** polarized:
with linear **patterns** that are radial/toroidal/helical

Scoring simulations with polarization: Image metrics

Unresolved linear polarization fraction

$$|m|_{\text{net}} = \frac{\sqrt{(\sum_i Q_i)^2 + (\sum_i U_i)^2}}{\sum_i I_i}$$

Unresolved circular polarization fraction (from ALMA)

$$|v|_{\text{net}} = \frac{|\sum_i V_i|}{\sum_i I_i}$$

Average resolved linear fraction

$$\langle |m| \rangle = \frac{\sum_i \sqrt{Q_i^2 + U_i^2}}{\sum_i I_i}$$

Average resolved circular fraction

$$\langle |v| \rangle = \frac{\sum_i |V_i/I_i|}{\sum_i I_i}$$

Azimuthal Linear structure 2nd mode (Palumbo+ 2020)

$$\beta_2 = \frac{1}{I_{\text{ring}}} \int_{\rho_{\text{min}}}^{\rho_{\text{max}}} \int_0^{2\pi} P(\rho, \varphi) e^{-2i\varphi} \rho d\varphi d\rho$$

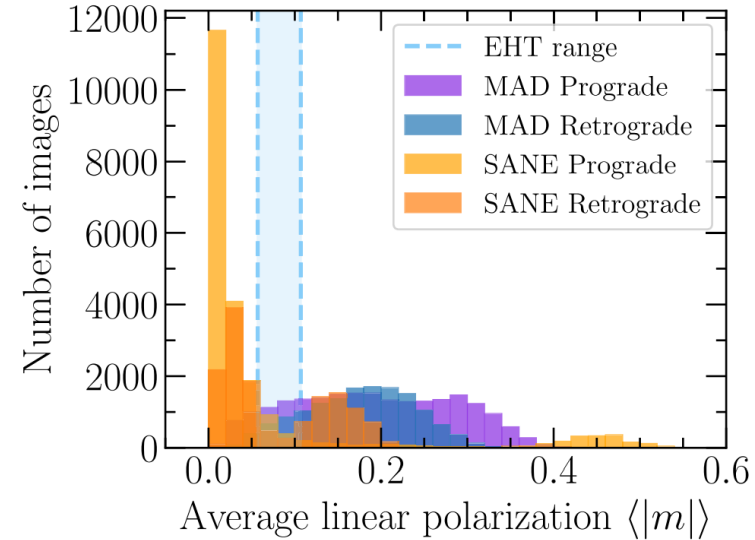
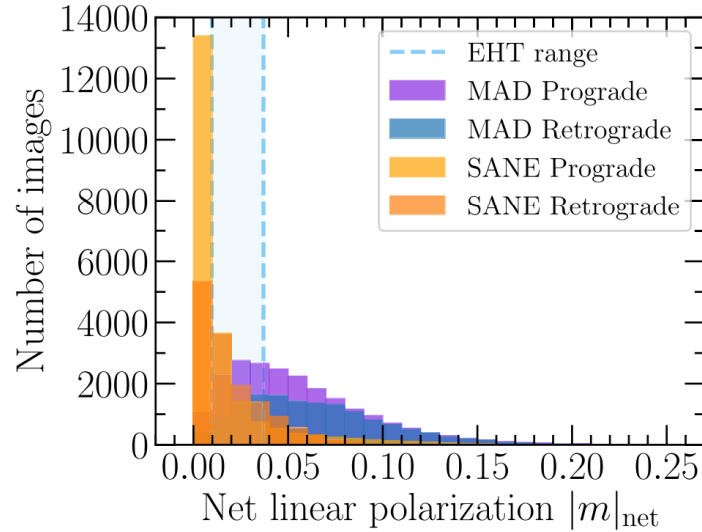
Table 3
Observational Constraints Applied to Our GRMHD Image Library

Parameter	Minimum	Maximum
m_{net}	1.0%	3.7%
v_{net}	-0.8%	0.8%
$\langle m \rangle$	5.7%	10.7%
$ \beta_2 $	0.04	0.07
$\angle \beta_2$	-163°	-129°
$\langle v \rangle$ (This Work)	0	3.7%

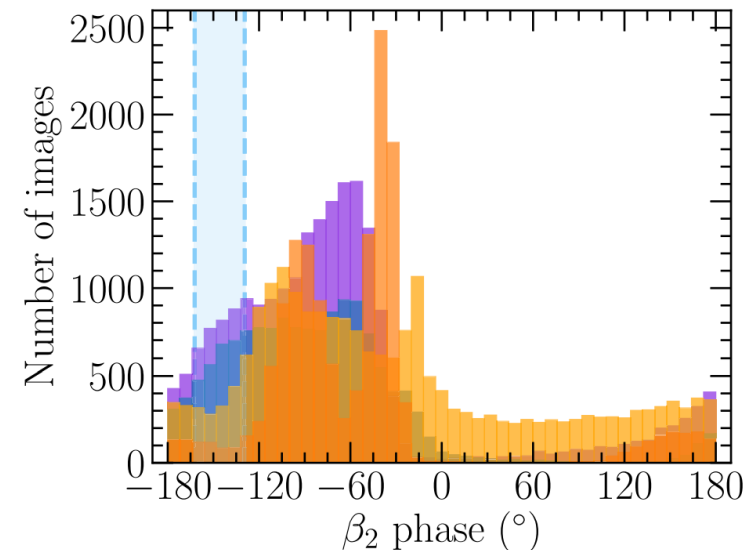
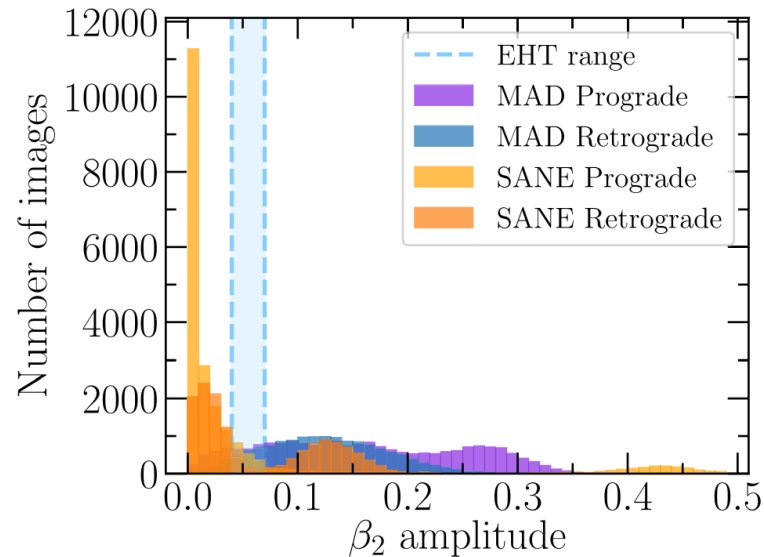
Note. Most of these constraints are inherited from Paper VII and were previously used to constrain models in Paper VIII. This work adds the new upper limit on $\langle |v| \rangle$.

Scoring simulations with linear polarization

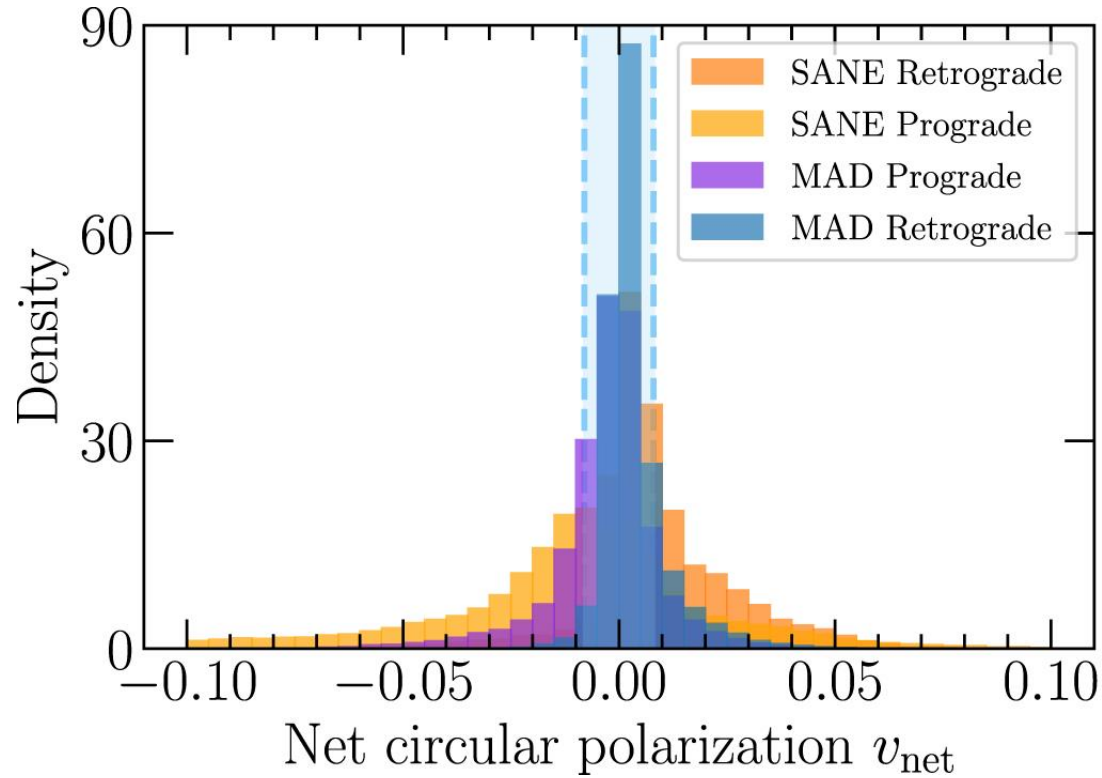
Unresolved and resolved linear polarization fraction:



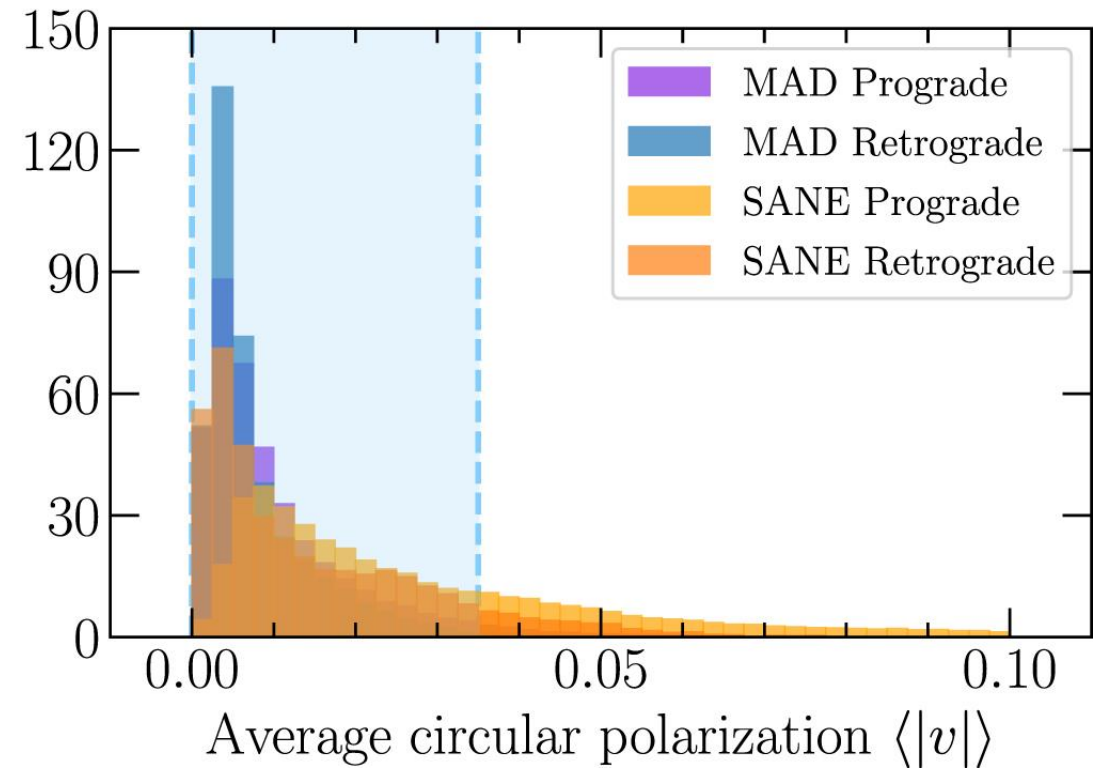
Azimuthal structure
2nd Fourier mode



GRMHD simulations naturally produce low circular polarization



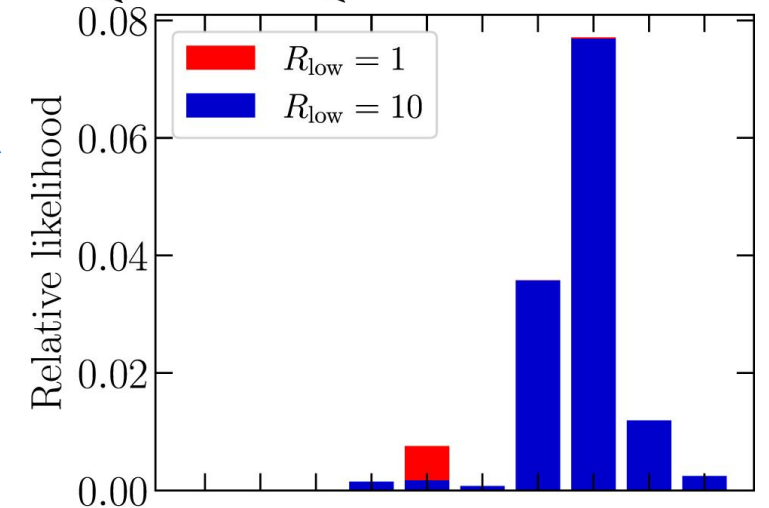
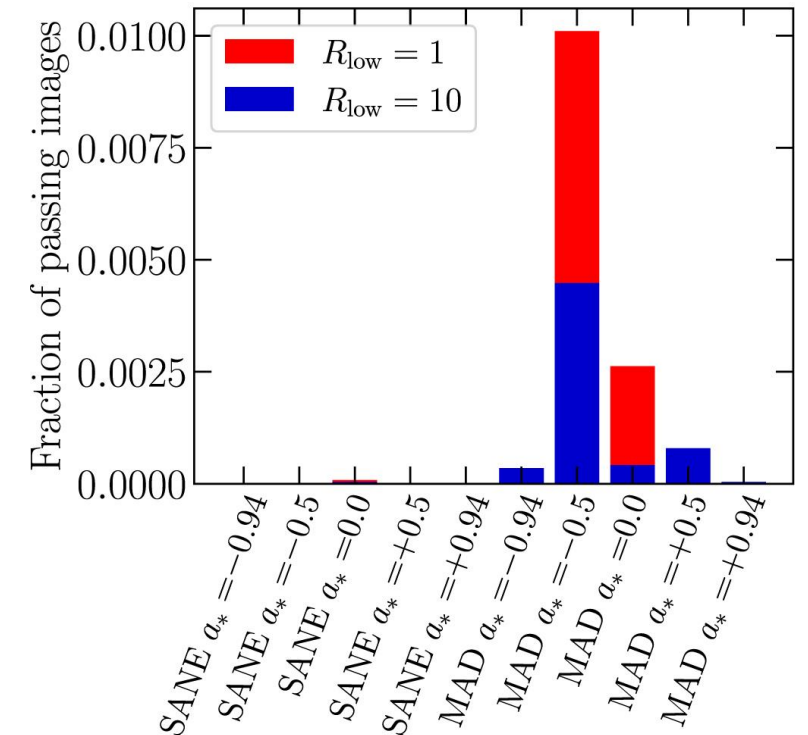
$$v_{\text{net}} = \frac{\int \mathcal{V} dA}{\int \mathcal{I} dA}$$



$$\langle |v| \rangle = \frac{\int |\mathcal{V}/\mathcal{I}| \mathcal{I} dA}{\int \mathcal{I} dA}$$

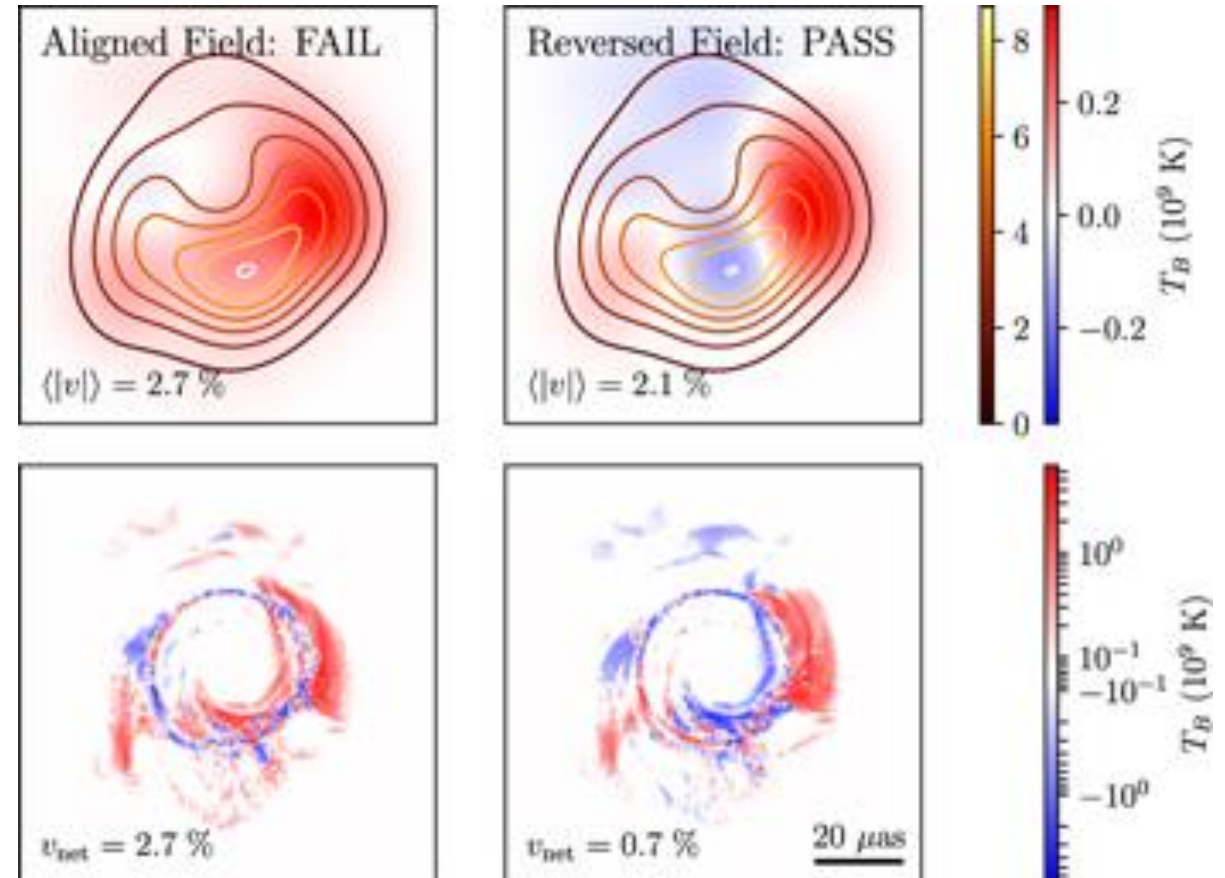
Polarimetric simulation scoring

- Two scoring approaches:
 - **'simultaneous'** (demand individual images satisfy all image constraints at once)
 - **'joint'** (compute a likelihood comparing distance between measured quantities and simulation mean with the simulation variance)
- **Both approaches strongly favor magnetically arrested (MAD) simulations**
- The two approaches differ in which electron heating parameters they favor.
- An additional constraint on the jet power rejects all surviving non-MAD simulations (and all spin-zero simulations)



Field orientation is very important!

- GRMHD is insensitive to the direction of the magnetic field, but polarized radiative transfer is **not**
- Changing the direction of the magnetic field changes:
 - sign of emitted V
 - direction of Faraday rotation
- images typically do not just flip sign when we reverse the field
 - circular polarization is typically produced via an interplay btw Faraday rotation and conversion,



Implications for M87*'s accretion

- Surviving models significantly tighten constraints on accretion rate from total intensity results:

$$\dot{M} \simeq (3 - 20) \times 10^{-4} M_{\odot} \text{ yr}^{-1}$$

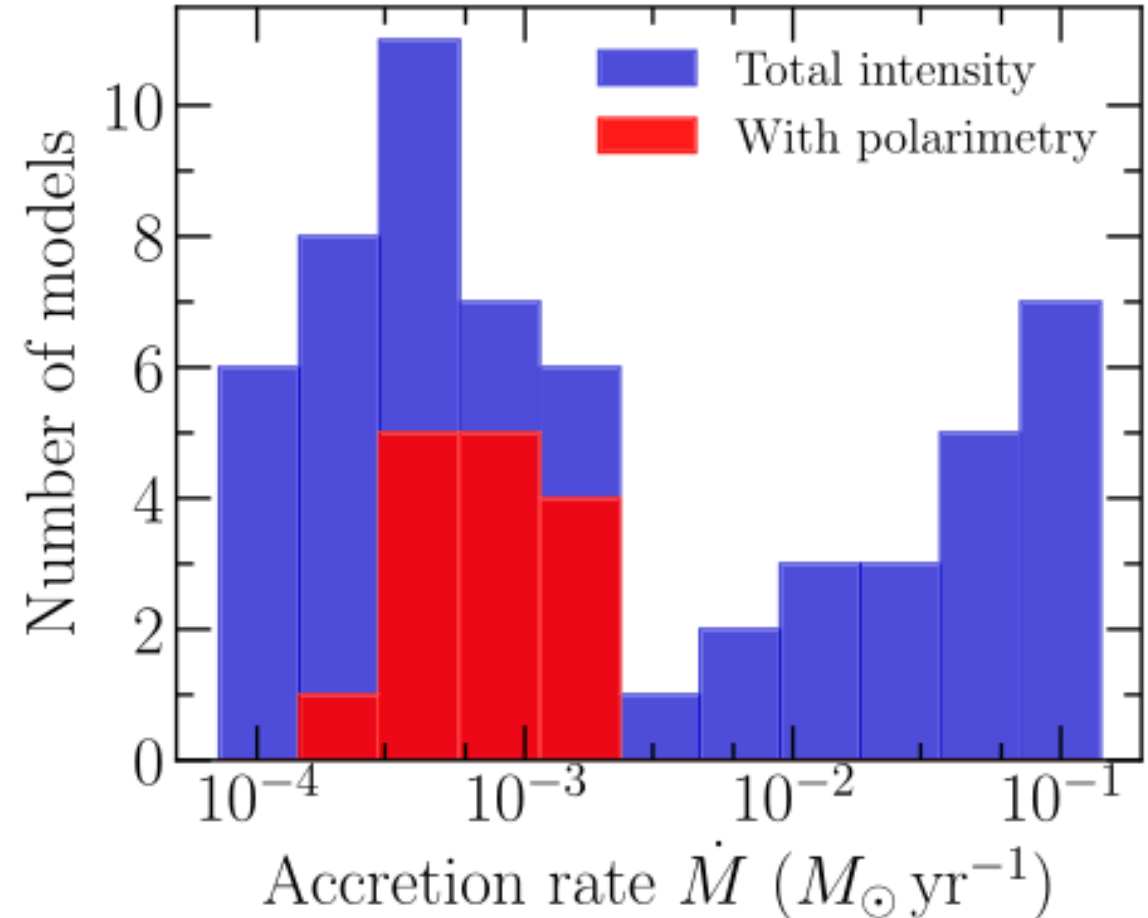
- Constrains the electron temperature, number density, and magnetic field strength (in agreement with estimates from simple one-zone models):

$$T_e \simeq (5 - 40) \times 10^{10} \text{ K}$$

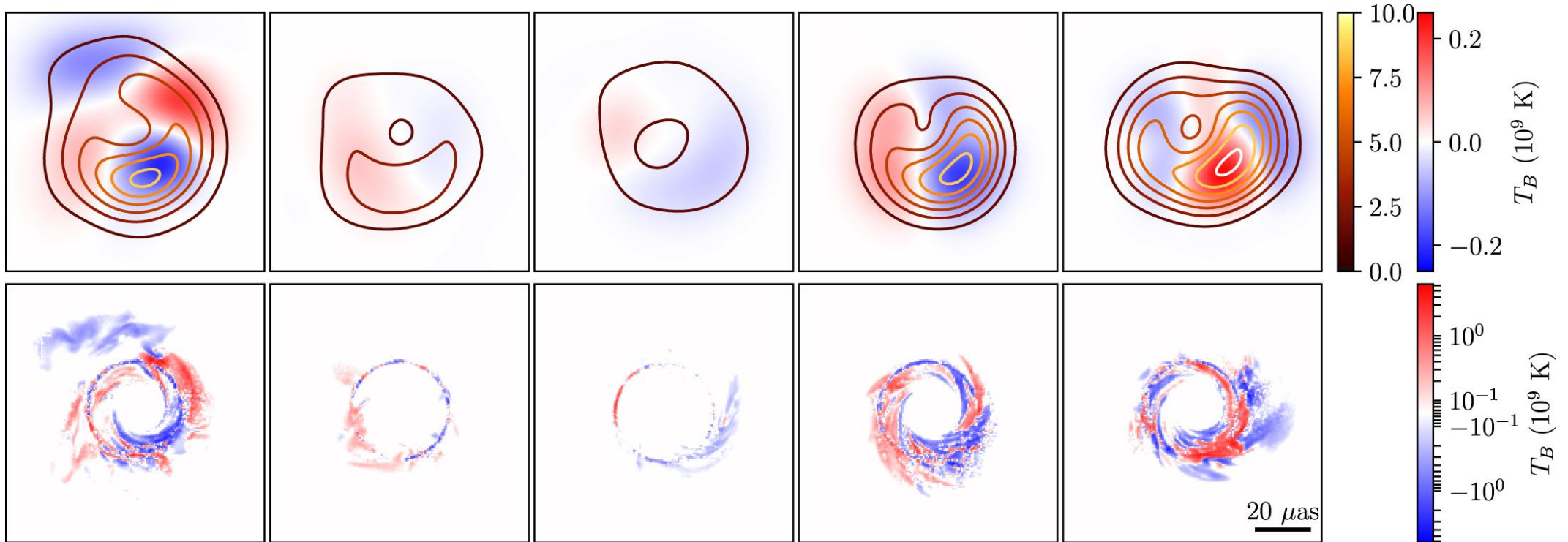
$$|B| \simeq (7 - 30) \text{ G}$$

$$n \sim 10^{4-5} \text{ cm}^{-3}$$

- Radiative efficiency $\sim 1\%$
 - Cooling is important!

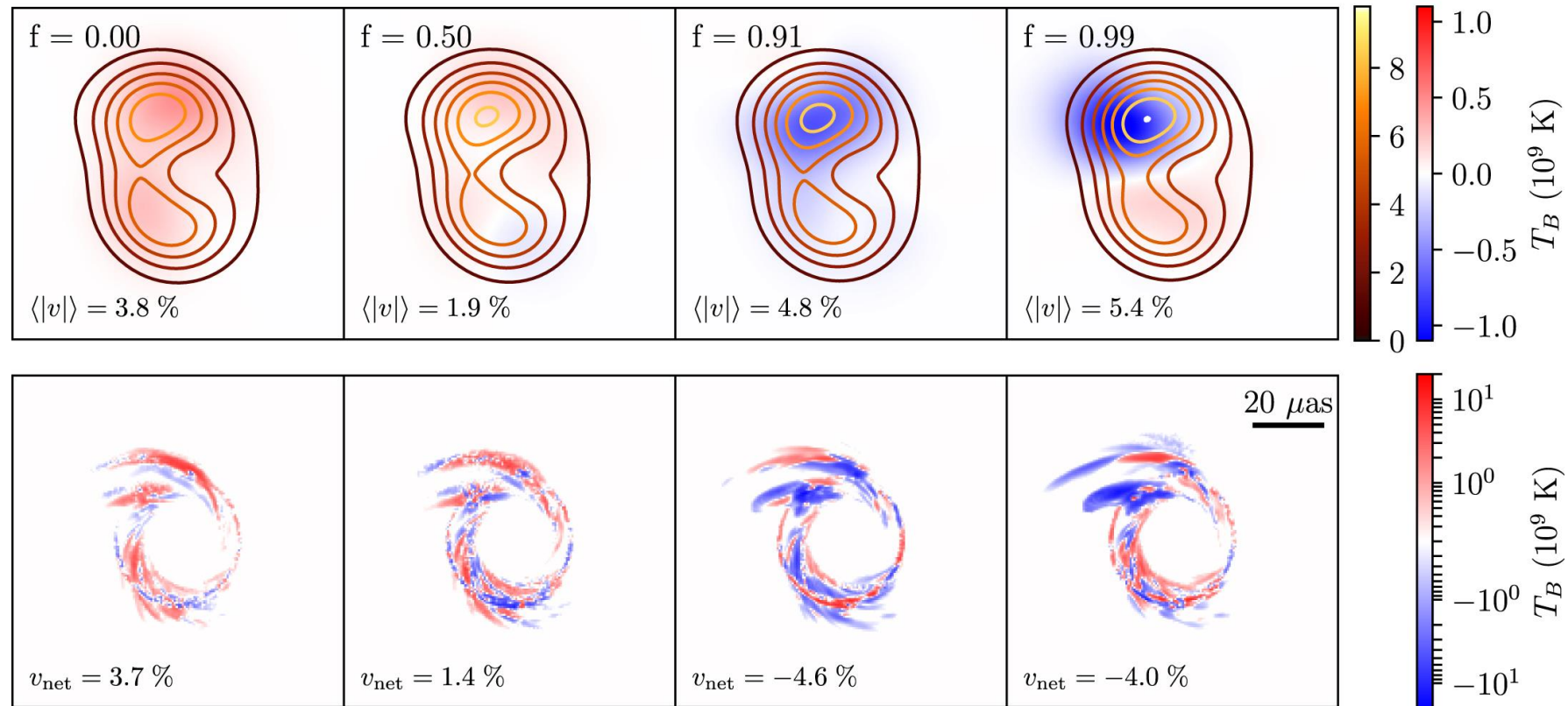


Passing simulations have diverse Stokes V morphologies



Detecting the Stokes V image structure with more sensitive observations will constrain our models further
Need more theoretical work to understand these morphologies!

Circular Polarization is sensitive to pairs,
but not in the way you might immediately think....



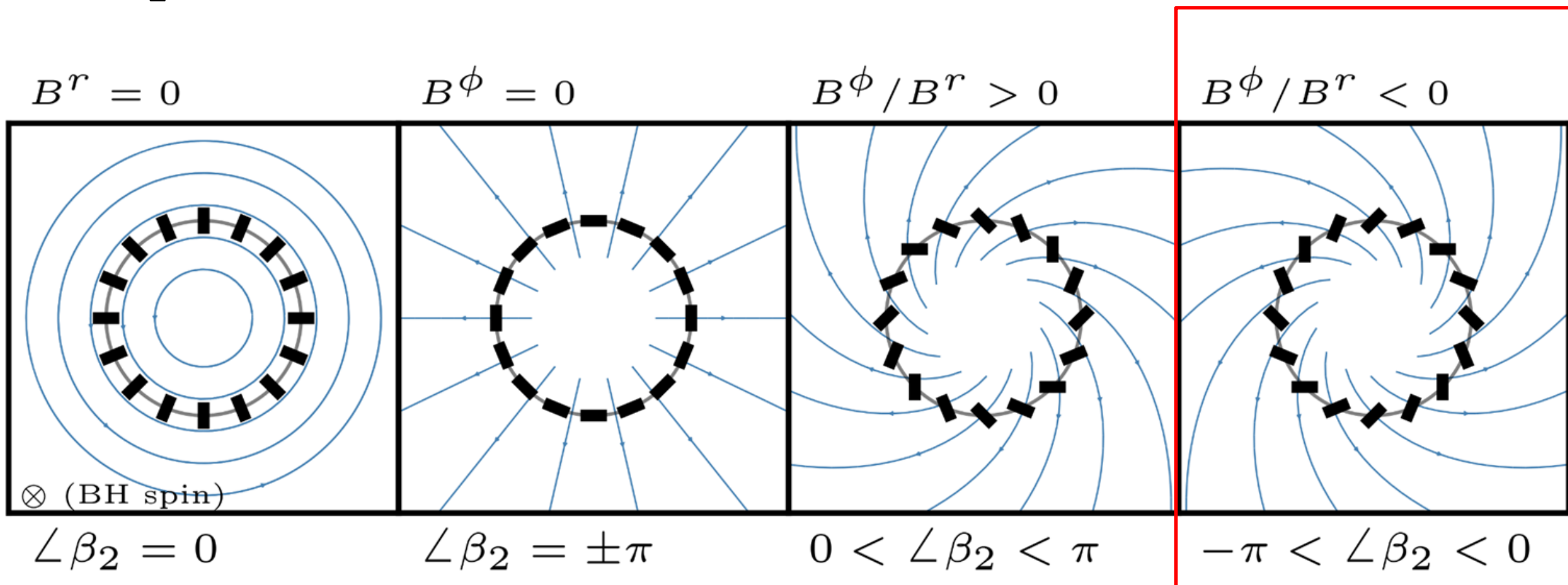
Conversion is the dominant source of Stokes V. It is **enhanced** by pairs
Faraday rotation is **reduced** by pairs
The interplay of these effects is complex

Connecting EHT images to electromagnetic energy flow

Chael, Lupsasca, Wong, Quataert 2023

[2307.06372](#)

$\arg(\beta_2)$ is connected to the ratio B^ϕ / B^r

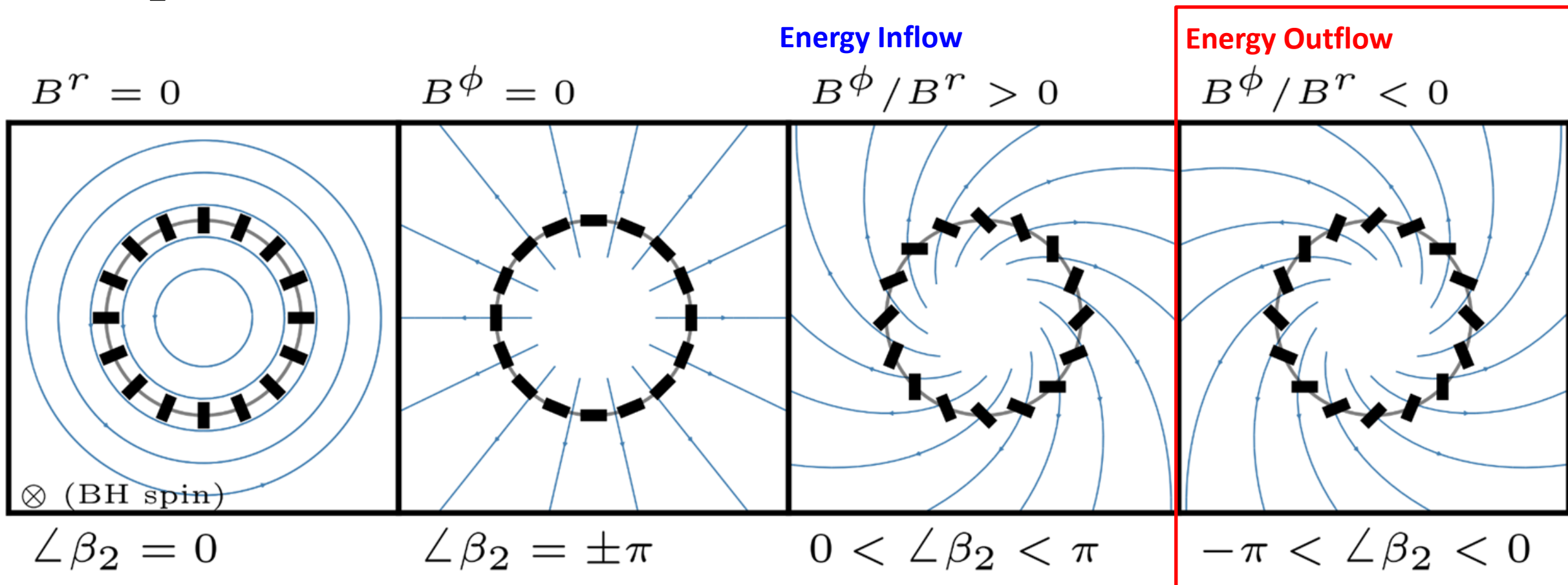


Cartoon picture:

- face on fields, no Faraday rotation, no optical depth, no relativistic parallel transport/abberation
- The BH spin is axis **into the screen** (EHT Paper V, 2019)

$$\angle \beta_2 \approx 2 \arctan \left(\frac{B^r}{r B^\phi} \right) \quad (\text{observer at } \theta_o = \pi)$$

$\arg(\beta_2)$ is connected to the **electromagnetic energy flux**



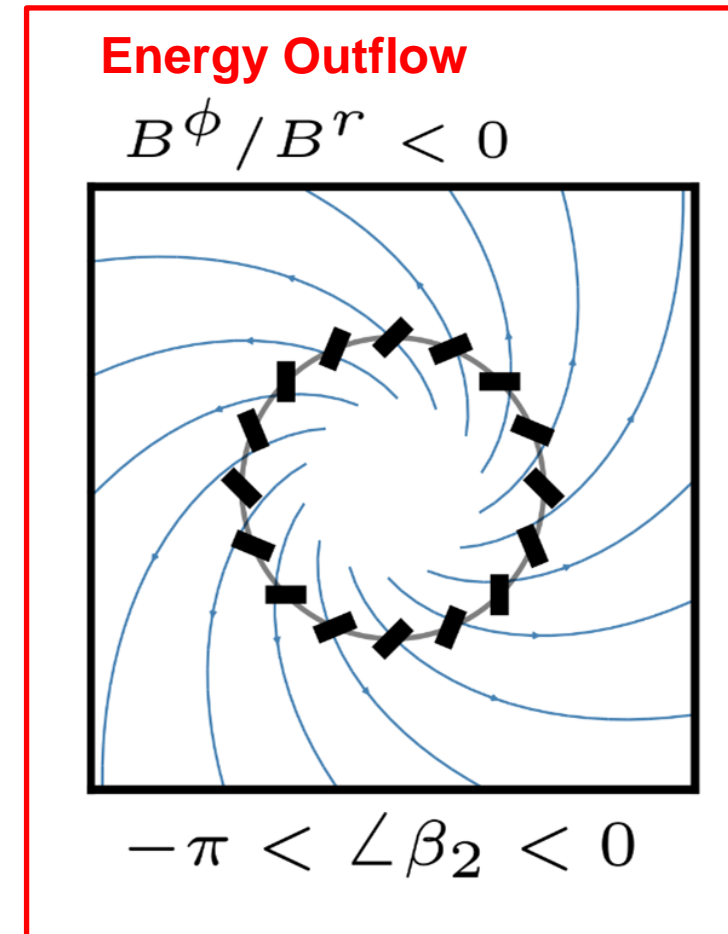
Poynting flux (Boyer-Lindquist Coordinates):

$$\mathcal{J}_\mathcal{E}^r = -T_{t\text{EM}}^r = -B^r B^\phi \Omega_F \Delta \sin^2 \theta.$$

↑ fieldline angular speed

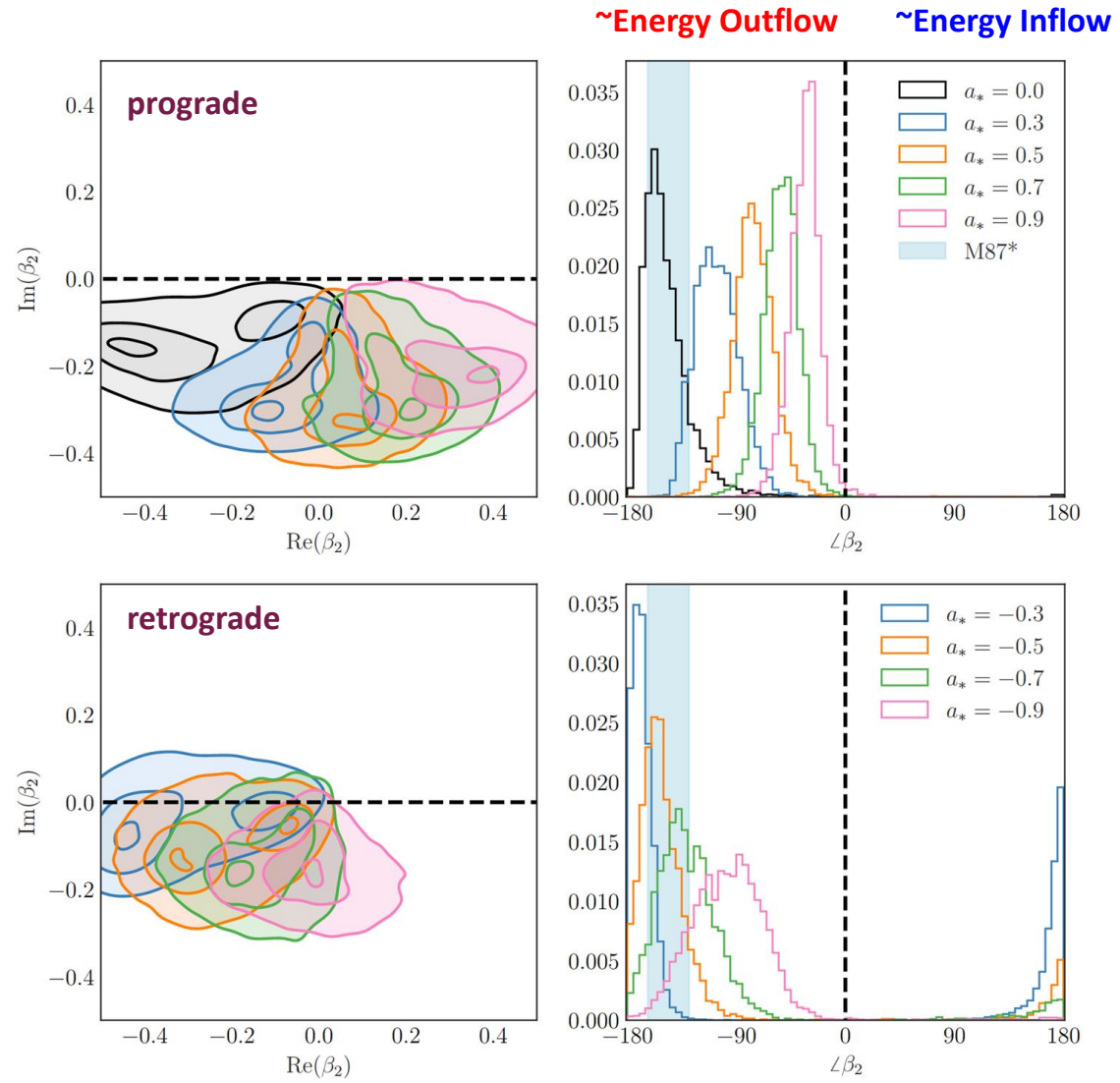
$\arg(\beta_2)$ is connected to the **electromagnetic energy flux**

- The sign of $\arg(\beta_2)$ is connected to the direction of Poynting flux
- Ignoring Faraday effects, The EHT's measurement $-163 \text{ deg} < \arg(\beta_2) < -129 \text{ deg}$ (Paper VII) implies **electromagnetic energy outflow in M87***
- This inference requires we know the **rotation direction**
 - We assume fieldlines **co-rotate** with the emitting material. (the angular velocity & spin vector is into the sky)
- Does this simple argument hold up in more complicated models of M87*?



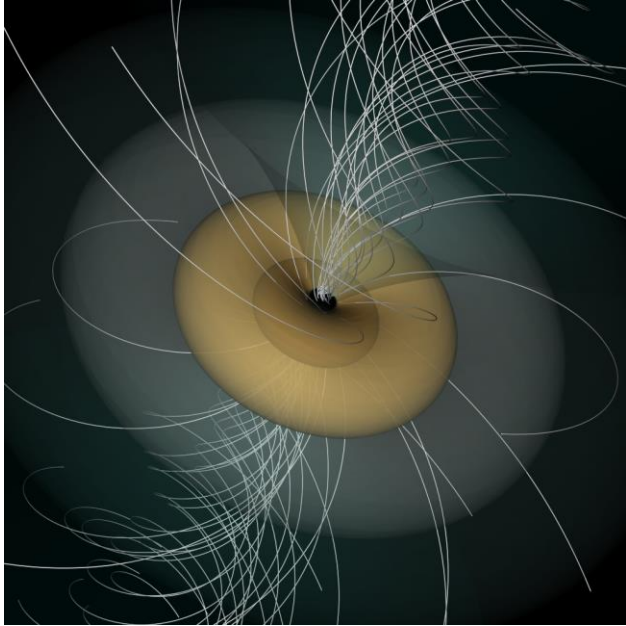
Does the relationship between $\arg(\beta_2)$ and energy flux persist in GRMHD models of M87*?

- M87* images from KORAL MAD simulations (Narayan+ 2022)
- 1600 snapshots covering different 6 different electron heating models (inclination fixed to M87* value)
- Almost all simulation images have $\arg(\beta_2)$ consistent with energy outflow in our simple picture
- $\arg(\beta_2)$ has the **same qualitative dependence on spin** as in the BZ monopole model

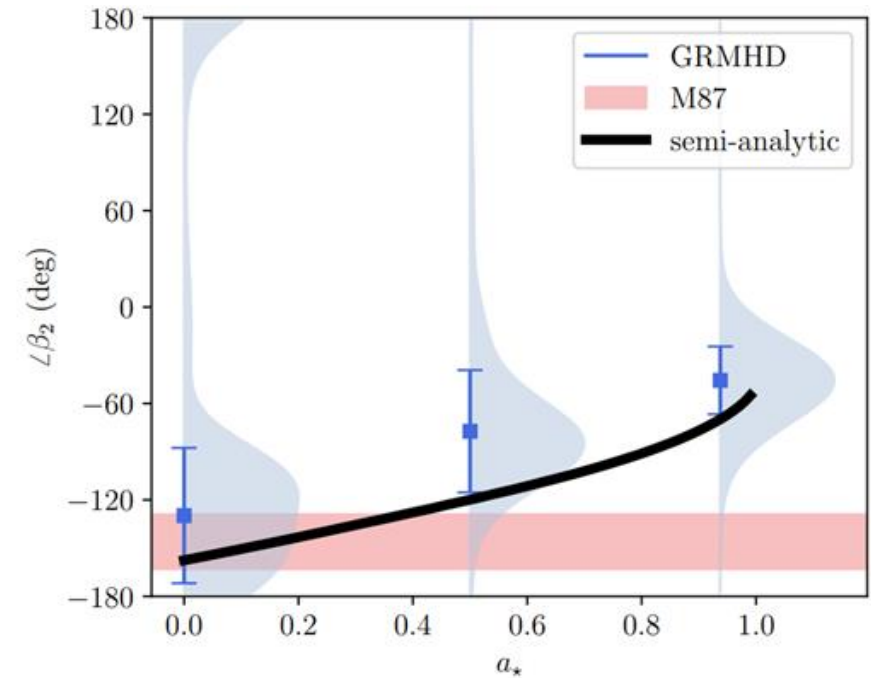
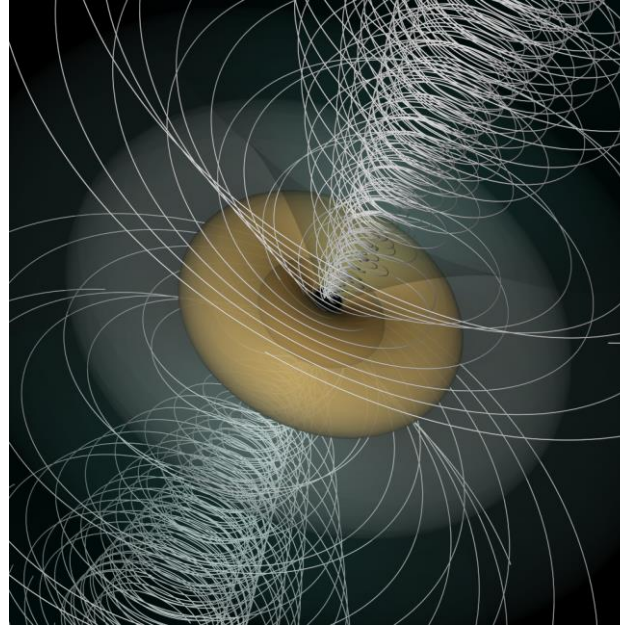


$\arg(\beta_2)$ has a strong dependence on BH spin in these models

Low Spin



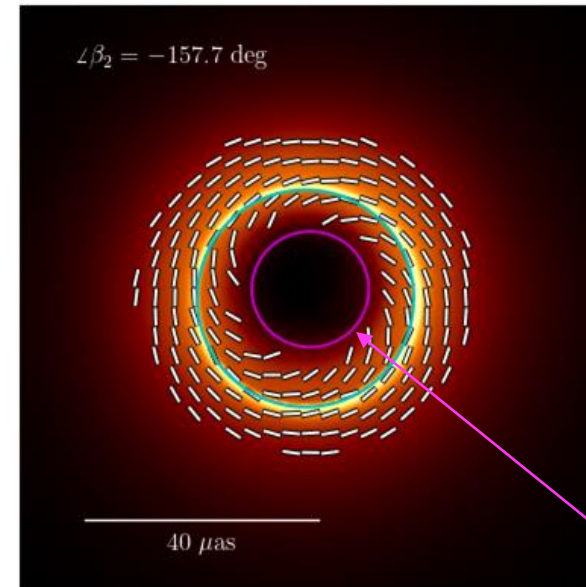
High Spin



- BH spin winds up initially radial fields, so that $B^\phi / B^r < 0$
- The field pitch angle increases with spin
- Increased field winding will
 - ☐ increase the Poynting flux (BZ jet power)
 - ☐ make the observed polarization more radial

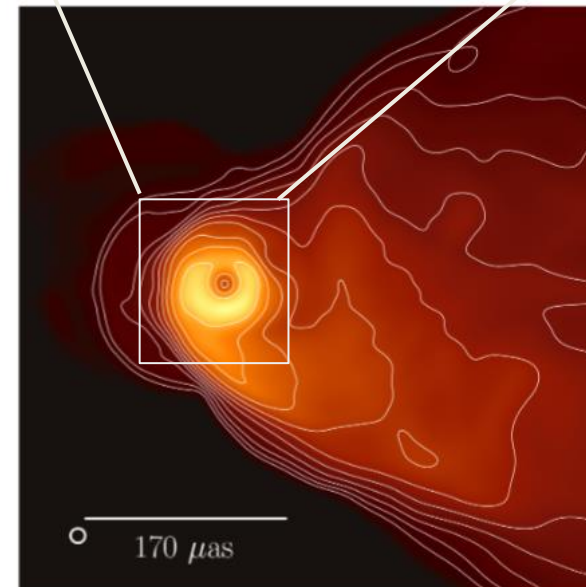
EHT/ngEHT next steps

- EHT Paper VII measurements of $\arg(\beta_2)$ suggest **electromagnetic outflow on scales of $\sim 5M$ in M87***.
- We can't yet be 100% sure if this energy outflow
 - is spin powered
 - or powers the large-scale jet
 - **the ngEHT could answer these questions!**
- We need **high-dynamic range, polarized ngEHT images** to:
 - Measure $\arg(\beta_2)$ **down to the horizon**
 - Connect the energy flux **from horizon scales out through the jet base**



Goal 1:
measure
energy flux
**down to
horizon**

“inner shadow”



Goal 2:
measure
energy flux
**out through
jet base**

Takeaways:

- The EHT has finally analyzed M87* in full polarization
- The structure of linear polarization is robustly constrained. Circular polarization is detected but the structure is not constrained.
- EHT linear polarization images show **~20% polarization** with an **azimuthal pattern** of polarization angles at 20 microarcsec scales. Circular polarization on these scales is **<4%**
- The EHT images can be used to constrain GRMHD simulation models of the emission region:
 - self-consistently including Faraday rotation and conversion effects is important
- The polarization data singles out magnetically arrested models:
 - **the magnetic field is dynamically important at the event horizon in M87***
 - These models naturally produce enough Faraday rotation to explain observed RM and low linear and circular polarization fractions
- The azimuthal structure of the linear polarization in M87* is consistent with outward Poynting flux
 - Simple model prediction is upheld in GRMHD simulation images.