Electron heating in simulations of Sgr A* (and M87)

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Ringberg, November 2, 2018









What does a black hole look like?

$$r_{\rm shadow} = \sqrt{27}GM/c^2$$

Black Hole Image Reconstruction with the EHT

(i.e. the other half of my work – ask me more later!)



Sagittarius A*

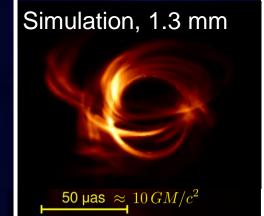
 $M_{BH} = (4.10 \pm 0.03) \times 10^6 M_{\odot}$

 $D = (8.12 \pm 0.03) \text{kpc}$

Gravity Collaboration, 2018 $d_{
m shadow}pprox 50\mu{
m as}$

GMVA, 3.5 mm (Issaoun+ 2018)

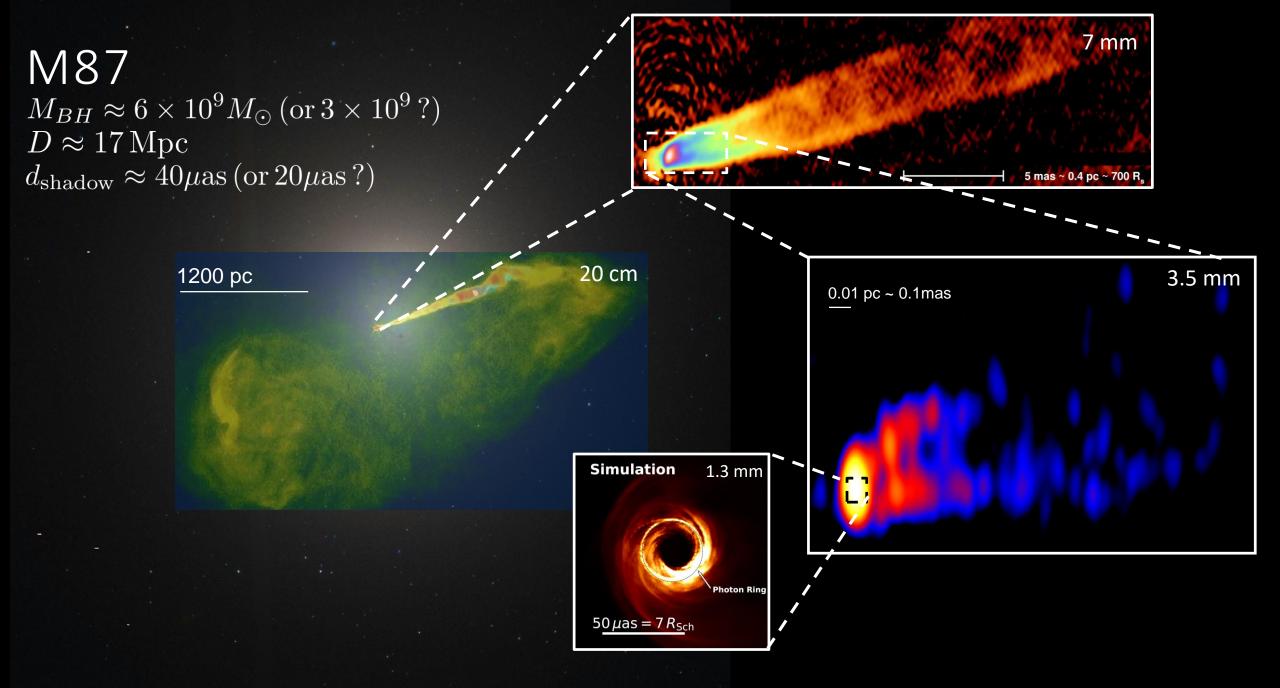
100 μ as $pprox 20 \, GM/c^2$



0.1°

Image credits: K.Y. Lo (VLA), UCLA Galactic Center Group (Keck), Sara Issaoun (GMVA+ALMA 3mm image)

 $\sim 10^6\,GM/c^2$



What will the EHT see?

- 1. Spacetime geometry
 - -The shadow of the black hole. Spin?
- 2. Fluid dynamics
 - -How is stuff moving? Jet/disk/outflow?
- 3. Electron (non)thermodynamics.
 - -Where are the emitting electrons?
 - -What is their distribution?

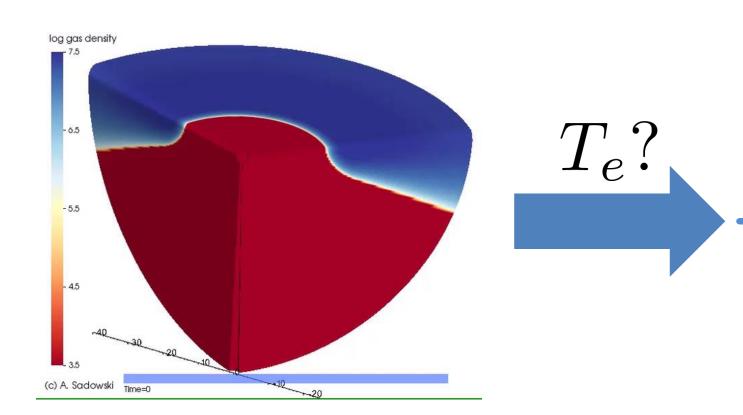
Sgr A* and M87 are Two-Temperature Accretion Flows

- Low densities in hot flows
 - → inefficient Coulomb coupling between ions and electrons.

Generally expect electrons to be cooler than ions.

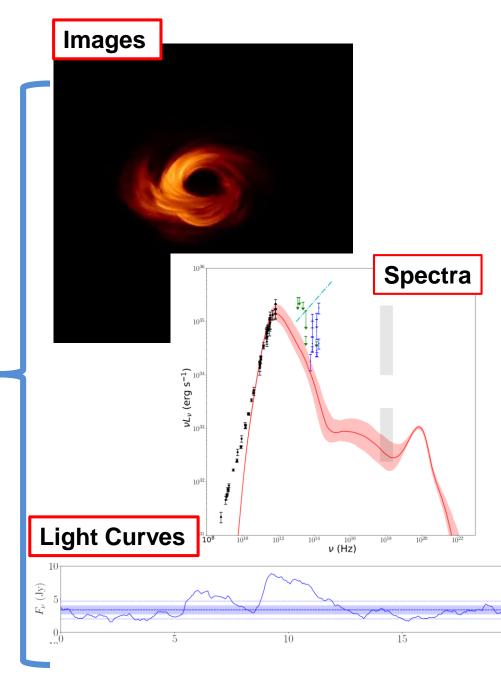
• But if electrons are **heated** much more, they can remain hotter.

From simulations to observables



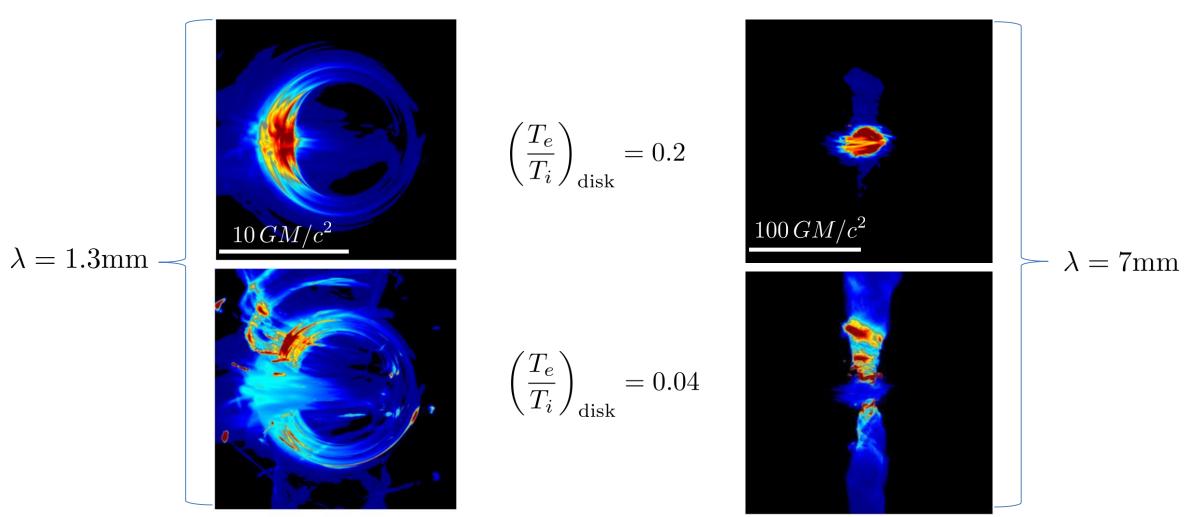
GRMHD Simulations

Usually evolve a single fluid and magnetic field



Fixed temperature ratios in postprocessing.

(Mościbrodzka+ 2014)



Different temperature ratios applied to the same simulation produce quite different images!

also: Dexter+2010, Moscibrodzka+ 2009, 2013, 2014, Chan+ 2015a, 2015b, Scherbakov+2012 ...

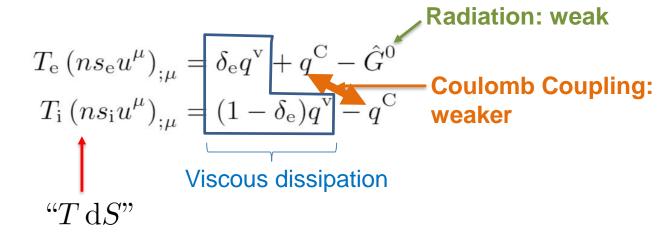
Goal: investigate the effects of microscale electron heating in **self-consistent** two-temperature simulations of the EHT targets Sgr A* and M87.

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-Using the code KORAL: (Sądowski+ 2013, 2015, 2017)
-See also previous work by:
Ressler+ 2017 (Sgr A*)
Ryan+ 2018 (M87)
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Two-Temperature GRRMHD Simulations

Total fluid quantities are evolved as in single-temperature GRRMHD

 Electron and ion energy densities are evolved via the 1st law of thermodynamics:

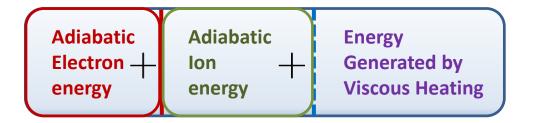


Adiabatic compression/expansion

Electron & Ion Heating

 The total dissipative heating in the simulation is internal energy of the total gas minus the energy of the components evolved adiabatically.

$$q^{v} = \frac{1}{\Delta \tau} \left[u - u_{i \text{ adiab}} - u_{e \text{ adiab}} \right]$$

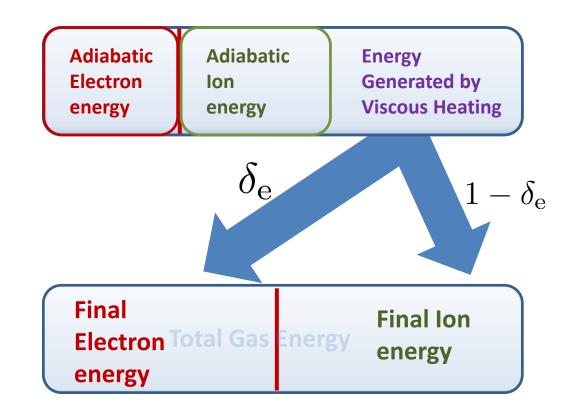


Total Gas Energy

Electron & Ion Heating

 The total dissipative heating in the simulation is internal energy of the total gas minus the energy of the components evolved adiabatically.

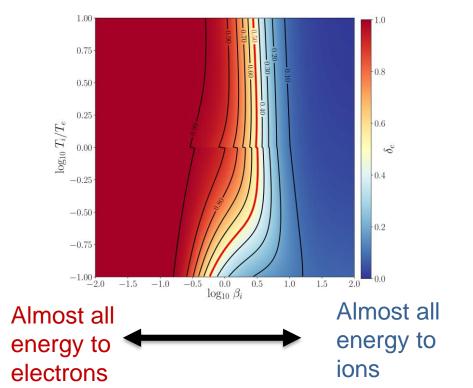
• **Sub-grid physics** must be used to determine what fraction of the dissipation goes into the electrons.



Sub-grid Heating Prescriptions

Landau-Damped Turbulent Cascade (Howes 2010)

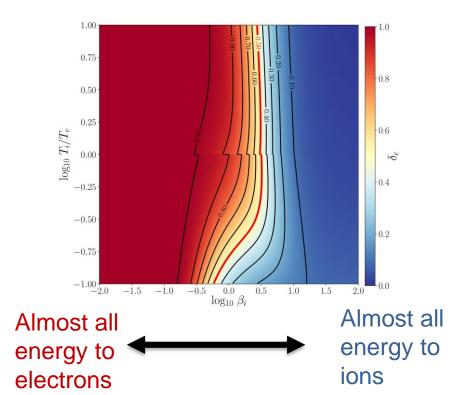
- Non-relativistic physics .
- Predominantly heats electrons when magnetic pressure is high, and vice versa



Sub-grid Heating Prescriptions

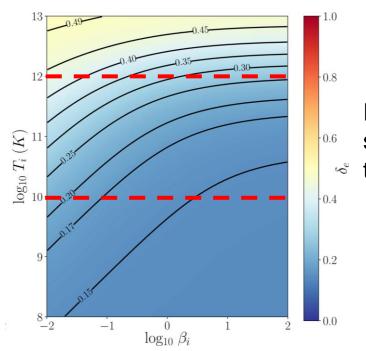
Landau-Damped Turbulent Cascade (Howes 2010)

- Non-relativistic physics .
- Predominantly heats electrons when magnetic pressure is high, and vice versa



Magnetic Reconnection (Rowan+ 2017)

- Based on PIC simulations of transrelativistic reconnection.
- Always puts more heat into ions
- Constant nonzero δ_e at low magnetization.



Relatively flat over significant range of temperature, beta.

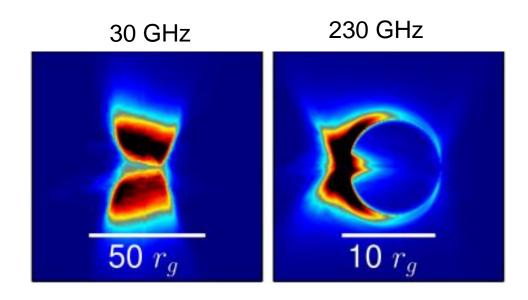
also: Kawazura+ 2018 (turbulent damping) Werner+ 2018 (reconnection)

Sgr A*

(Chael+ 2018b, arXiv: 1804.06416)

Previous work: Ressler et al. 2017

- A 3D, two-temperature simulation with relatively high magnetic flux and using the turbulent cascade prescription.
- Natural disk-jet structure.
- Q: Is this structure dependent on electron heating & B field strength?



Our Sgr A* Simulations

- Four 3D simulations using KORAL
- one for each heating prescription at low (0) and high (0.9375) BH spins.

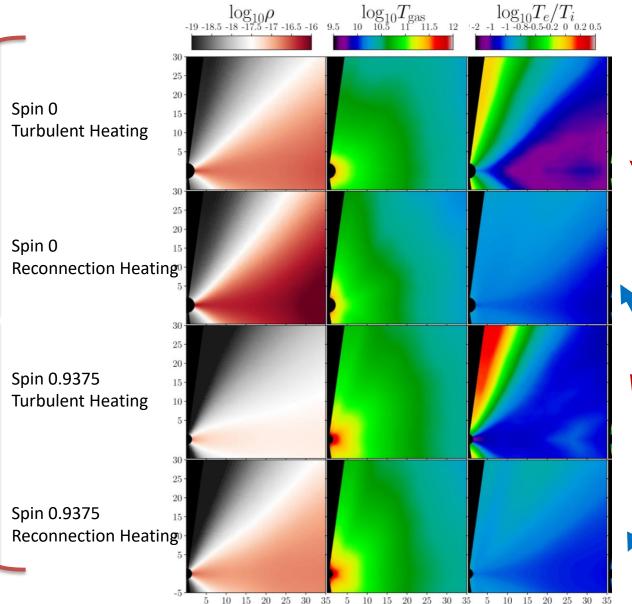
| Model | Spin | Heating | $\dot{M}(\dot{M}_{ m Edd})$ | $\Phi_{ m BH} \left((\dot{M}c)^{1/2} r_{ m g} \right)$ |
|----------------------|--------|-------------------|-----------------------------|---|
| H-Lo | 0 | Turb. Cascade | 3×10^{-7} | 5 |
| R-Lo | 0 | Mag. Reconnection | 7×10^{-7} | 4 |
| H-Lo R-Lo H-Hi | 0.9375 | Turb. Cascade | 2×10^{-7} | 6 |
| | 0.9375 | Mag. Reconnection | 3×10^{-7} | 3 |

Very **low** "MAD parameter" ~50 is saturation value for a Magnetically Arrested Disk

• Density is scaled to match 3.5 Jy at 230 GHz (Bower+ 2015).

Sgr A*: Temperature ratio

All are thick disks: density lower at high spin



Temperature ratio is highly stratified with polar angle for turbulent heating Electrons are hotter than ions in the jet

Relatively constant temperature ratio for reconnection Electrons are cooler everywhere

230 GHz mm movies

Spin 0 Turbulent Heating





Spin 0.9375 Turbulent Heating

Spin 0 Reconnection Heating

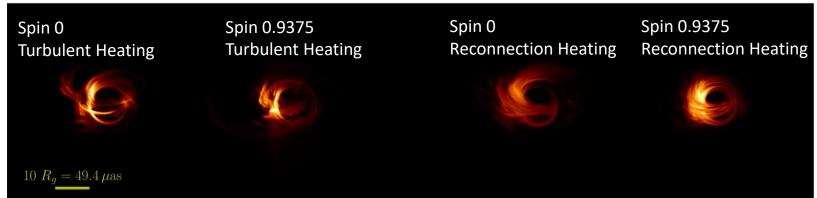




Spin 0.9375 Reconnection Heating

Image structure with frequency

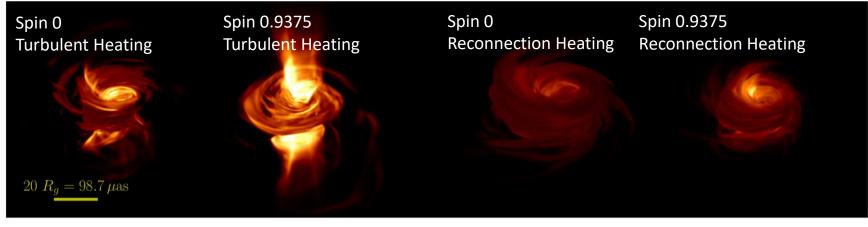
230 GHz



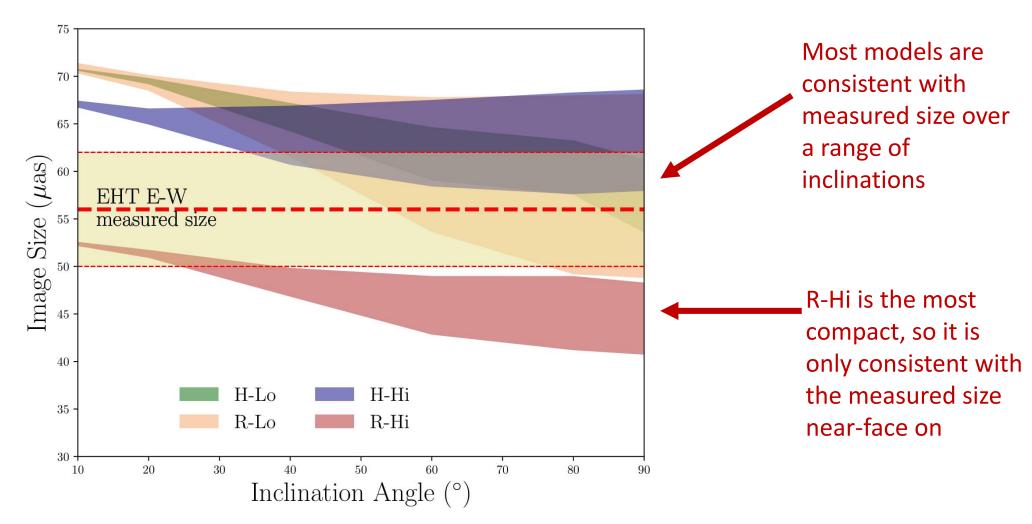
Where the EHT observes at 230 GHz, both heating prescriptions produce images with distinct black hole shadows

Turbulent heating makes lower frequency images anisotropic and jet dominated — exceeding estimates of intrinsic anisotropy when viewed at high inclination (Johnson+ 2018, Issaoun+ 2018)

43 GHz

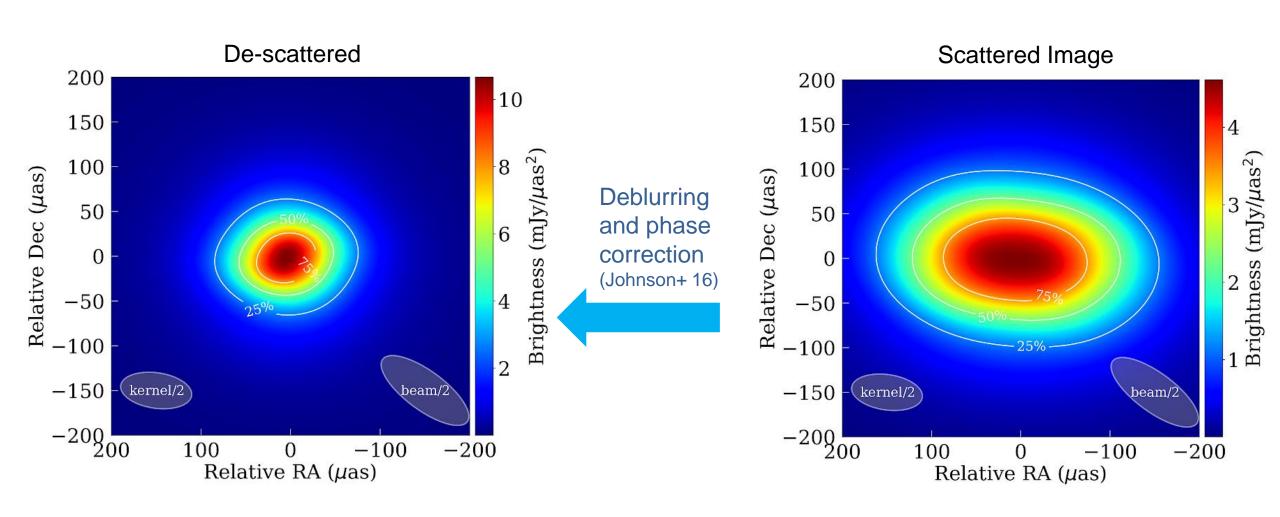


EHT 230 GHz size



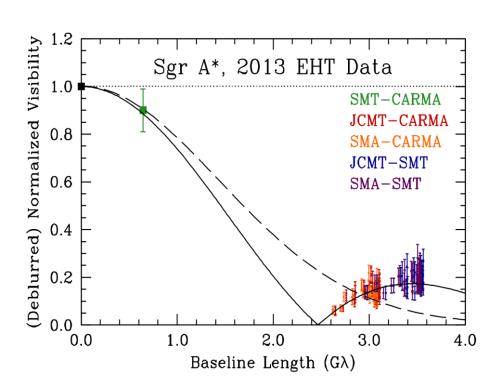
First Intrinsic Image of Sgr A* at 3.5 mm

and the first VLBI with ALMA (Issaoun+ 2018)

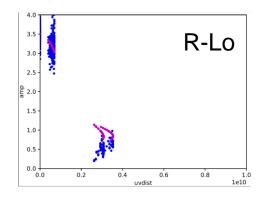


New constraints on Sgr A* asymmetry at 3.5 mm rule out edge-on jet!

Comparison with EHT 2013 visibilities.

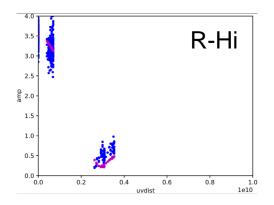


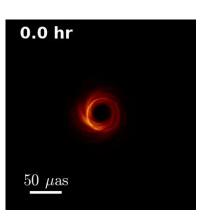
60 degree inclination – no visibility null



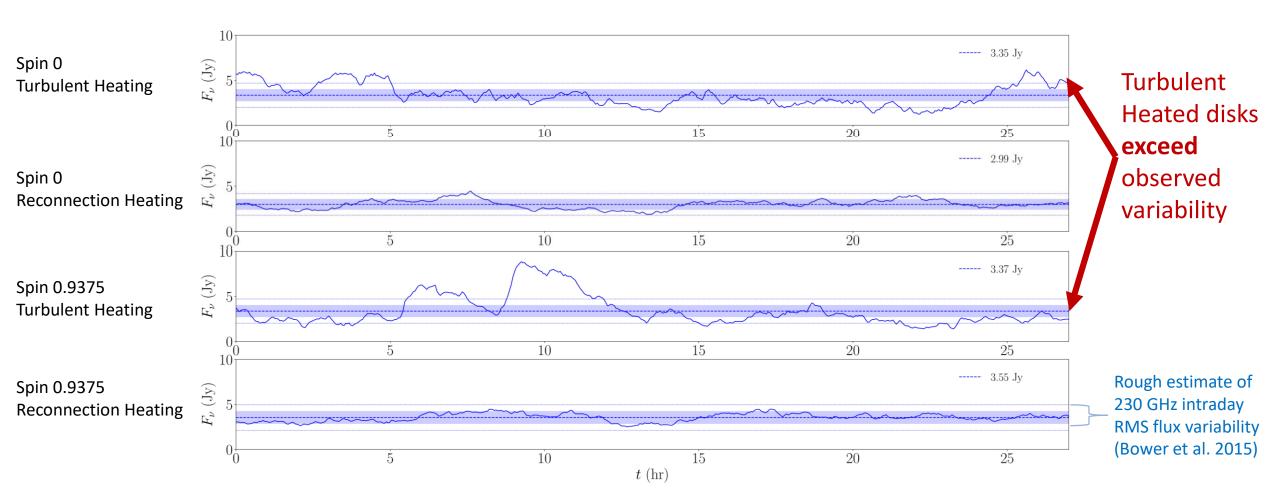


10 degree inclination – visibility null from symmetric ring

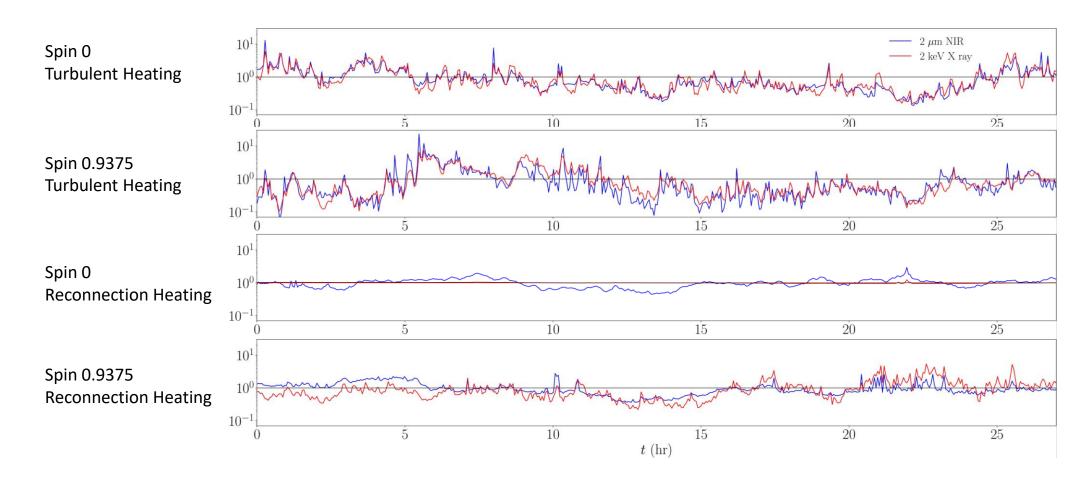




Sgr A*: 230 GHz variability

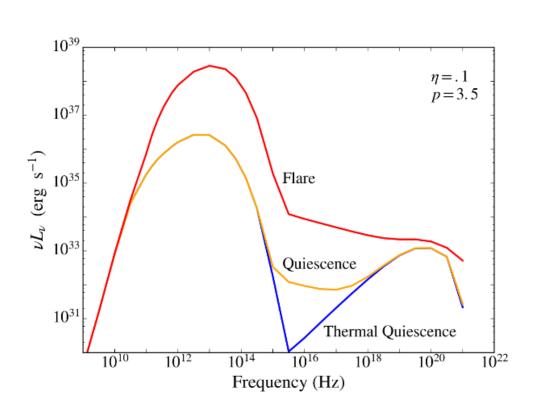


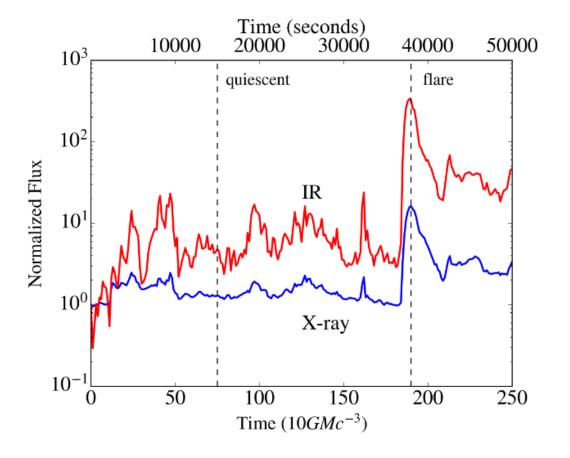
IR and X-ray variability: no flares



No models reproduce strong IR and X-ray flares → Nonthermal Electrons

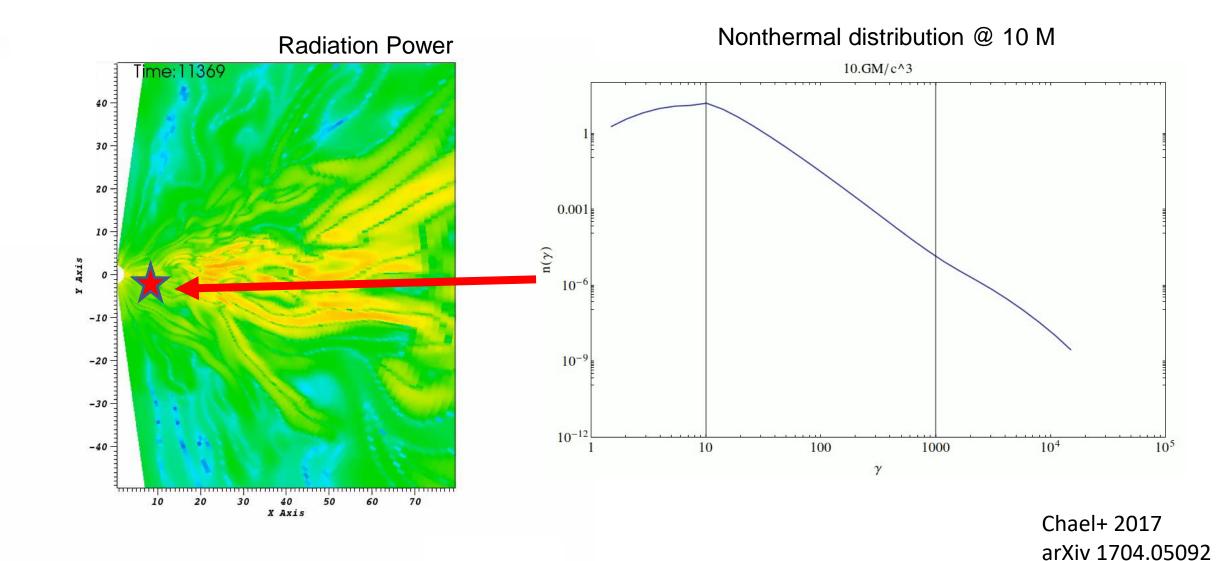
Nonthermal distributions added in postprocessing can produce correlated NIR & X-ray flares





Evolving nonthermal electrons in simulations

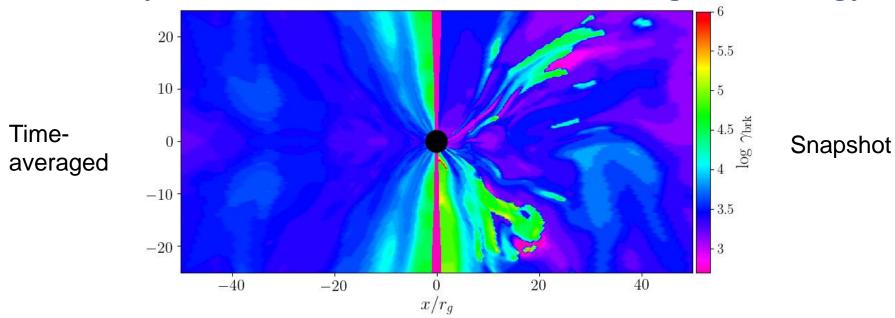
(Chael+ 2017)



Evolving **nonthermal** electrons in simulations (Chael+ 2017)

New method to self-consistently evolve non-thermal spectra in parallel with two-temperature fluid.

Spatial distribution of nonthermal cooling break energy



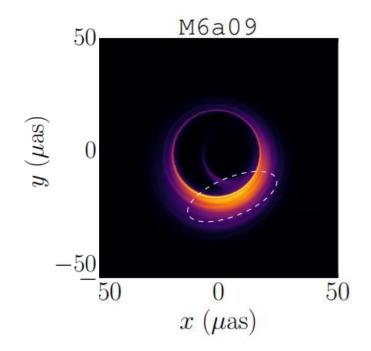
- First 3D simulations with realistic electron acceleration (Ball+ 2018) coming soon!
 - $_{ullet}$ Ball+ 2018: Magnetic reconnection PIC simulations give p=2.5 (Ponti+ 2017) at $\sigmapprox1$
 - Jet sheath as acceleration site?

M87

(Chael+ 2018b, arXiv: 1810.01983)

Previous work: Ryan et al. 2018

- 2D, two-temperature simulations with weak magnetic flux and using the turbulent cascade prescription at 2 BH masses.
- Good agreement with previous EHT measurements of image size for high mass case $(6 \times 10^9 M_{\odot})$.
- Jet power relatively weak, jet angle is narrow.

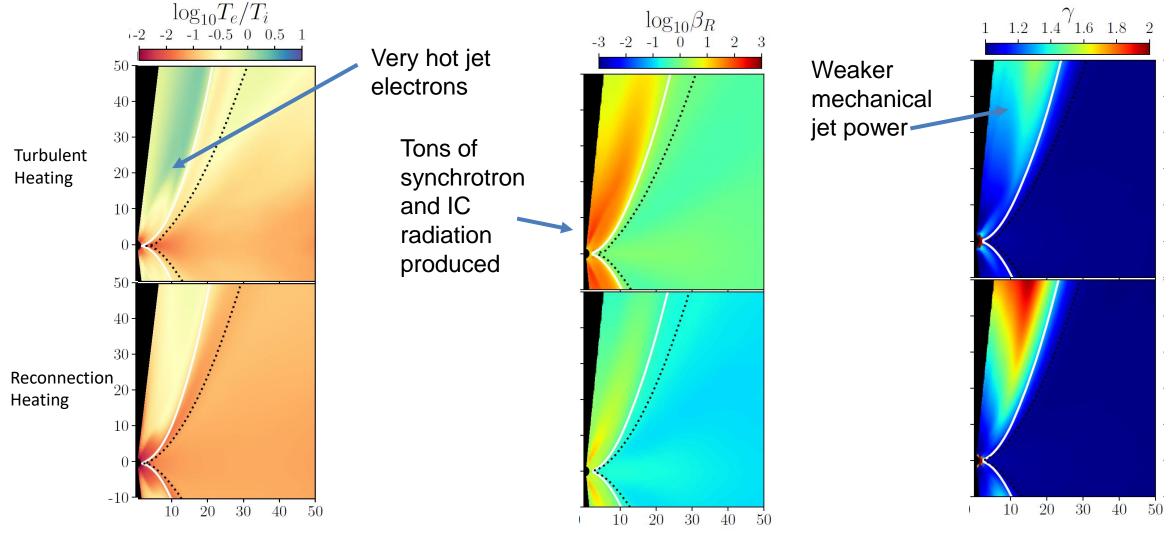


Our M87 Simulations

| Model | Spin | Heating | $\langle \dot{M}/\dot{M}_{ m Edd} angle$ | $\langle \Phi_{ m BH}/(\dot{M}c)^{1/2}r_{ m g} \rangle$ | $\langle P_{J(100)} \rangle \ [{\rm erg \ s^{-1}}]$ |
|-------|--------|-------------------|---|---|---|
| H10 | 0.9375 | Turb. Cascade | 3.5×10^{-6} | 54 | 6.6×10^{42} |
| ★ R17 | 0.9375 | Mag. Reconnection | 2.3×10^{-6} | 63 | 1.2×10^{43} |
| · | | | • | | |
| | | | | "MAD parameter" | Jet mechanical power |

- Both simulations are MAD.
- Density is scaled to match 0.98 Jy at 230 GHz.
- The mechanical jet power in R17 is in the measured range of $10^{43}-10^{44}$ erg/s.

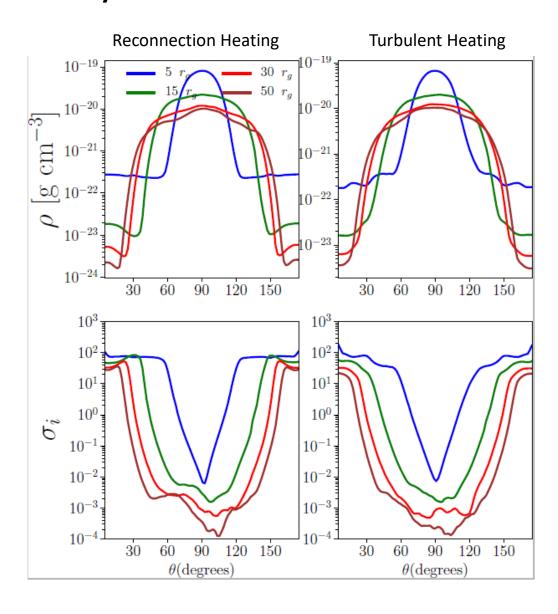
Electron Heating -> Jet Dynamics



Turbulent heating produces too much radiation at the jet base, which saps the jet power

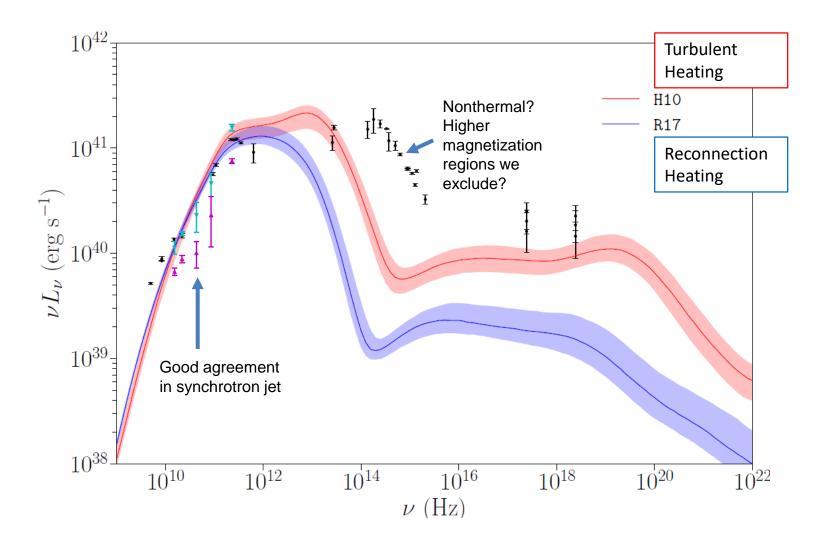
Electron Heating + Radiation → Dynamics!

Density floors: σ_i cut in radiative transfer



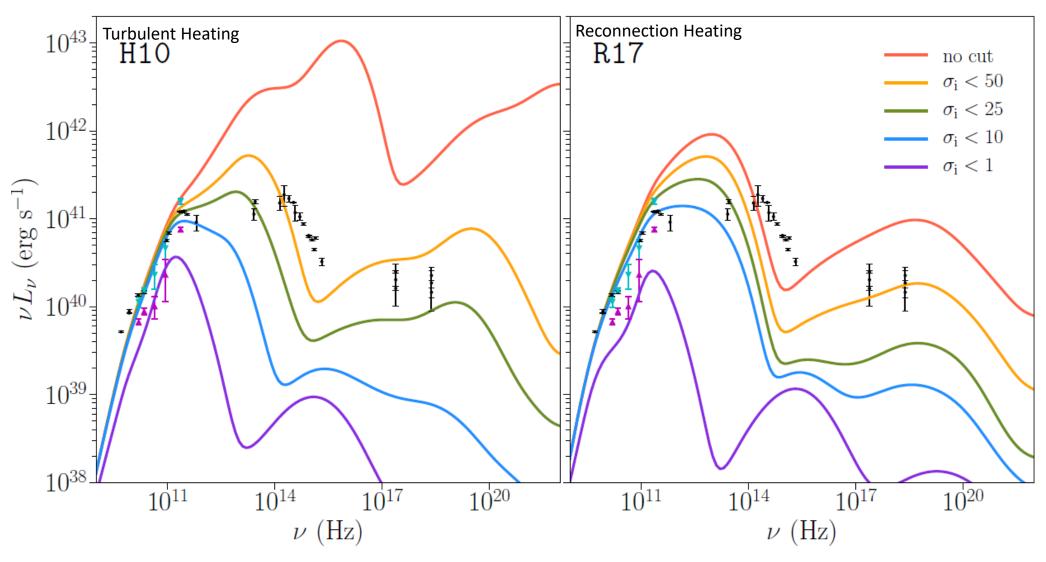
- Density floors are imposed in the simulation inner jet where $\sigma_{\rm i} \geq 100$
- We don't trust radiation from these regions, so when raytracing we only include regions where $\sigma_{\rm i} \leq 25$
- Spectra and images at frequencies ≥ 230 GHz depend strongly on the choice of cut!

M87 Spectrum



Data from Prieto+16
New points (cyan and magenta) from Akiyama+15, Doeleman+12, Walker+18, Kim+18, and MOJAVE

M87 Spectra: dependence on $\sigma_{\rm i}$ cut



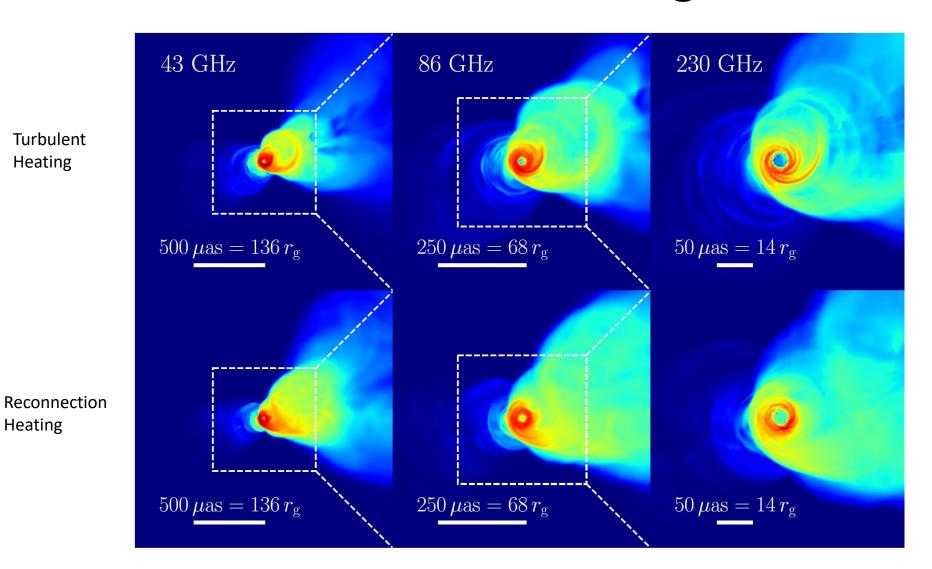
Data from Prieto+16
New points (cyan and magenta) from Akiyama+15, Doeleman+12, Walker+18, Kim+18, and MOJAVE

M87 Jets at millimeter wavelengths

Turbulent

Heating

Heating



Inclination angle (down from pole)

 17°

Disk/Jet rotation sense

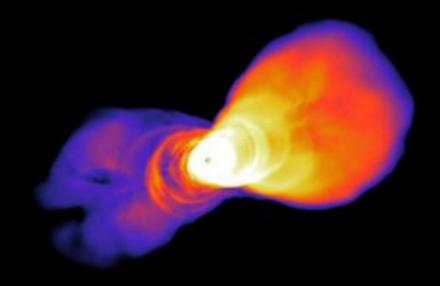


(Values from Walker+ 18)

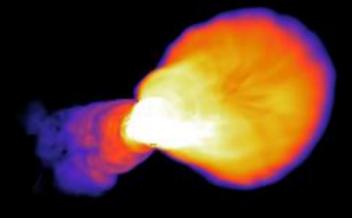
Wide apparent opening angles get larger with increasing frequency

43 GHz jets

0.0 yrTurbulent Heating



Reconnection Heating



 $500 \ \mu as$

43 GHz images – comparison with VLBA

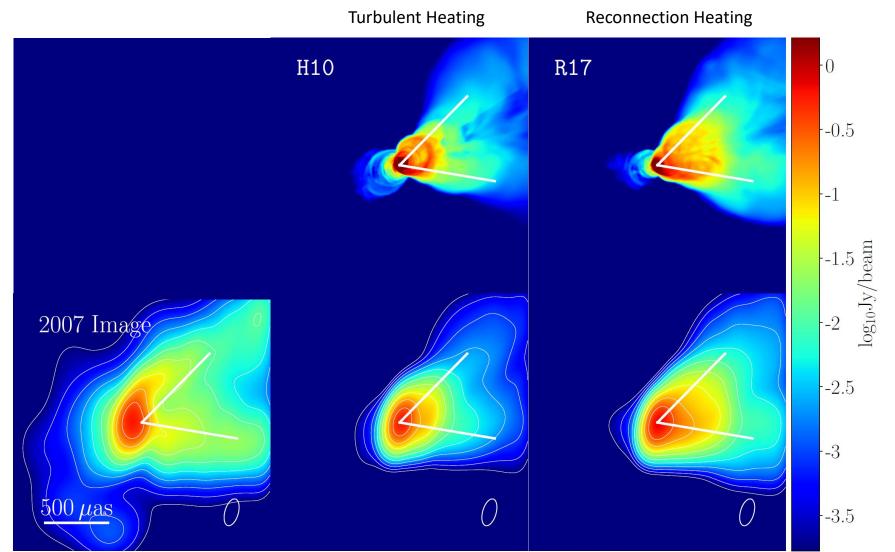
Walker+ 2018

High

VLBA

Resolution

Resolution

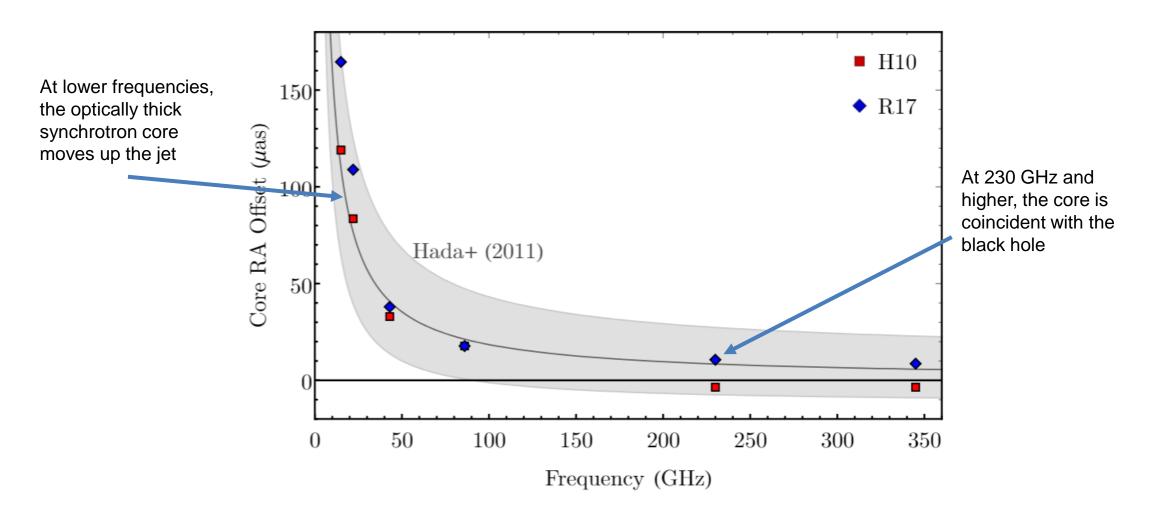


Apparent opening angle at 43 GHz:

55°

(Walker+ 2018)

M87 Core Shift



Agreement with measured core shift down to cm wavelengths.

What will M87 look like to the EHT at 230 GHz?

0.0 yr
Turbulent Heating

Reconnection Heating





What will M87 look like to the EHT at 230 GHz?

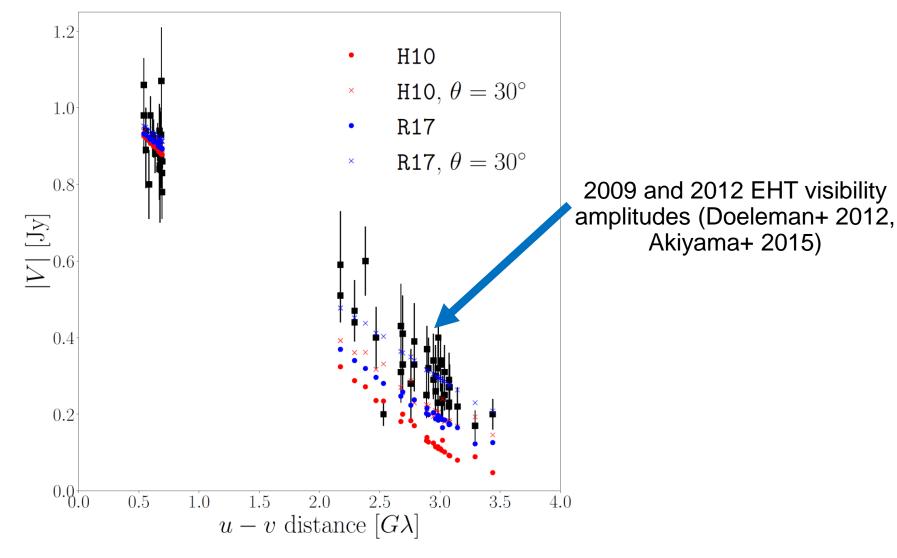
0.0 yr
Turbulent Heating

Reconnection Heating



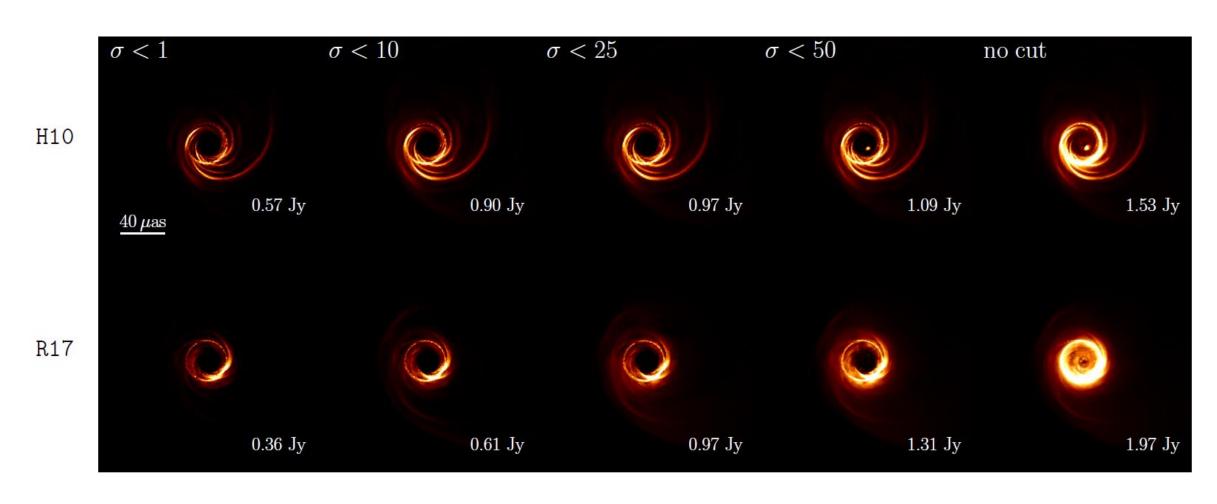


Current 230 GHz images are too big!



Changing the inclination **or** including more emission from magnetized regions near axis makes the emission more compact.

230 GHz images – dependence on σ_{i} cut



The image becomes more compact & counterjet dominated when we include more high-magnetization emission from the jet base!

Takeaways

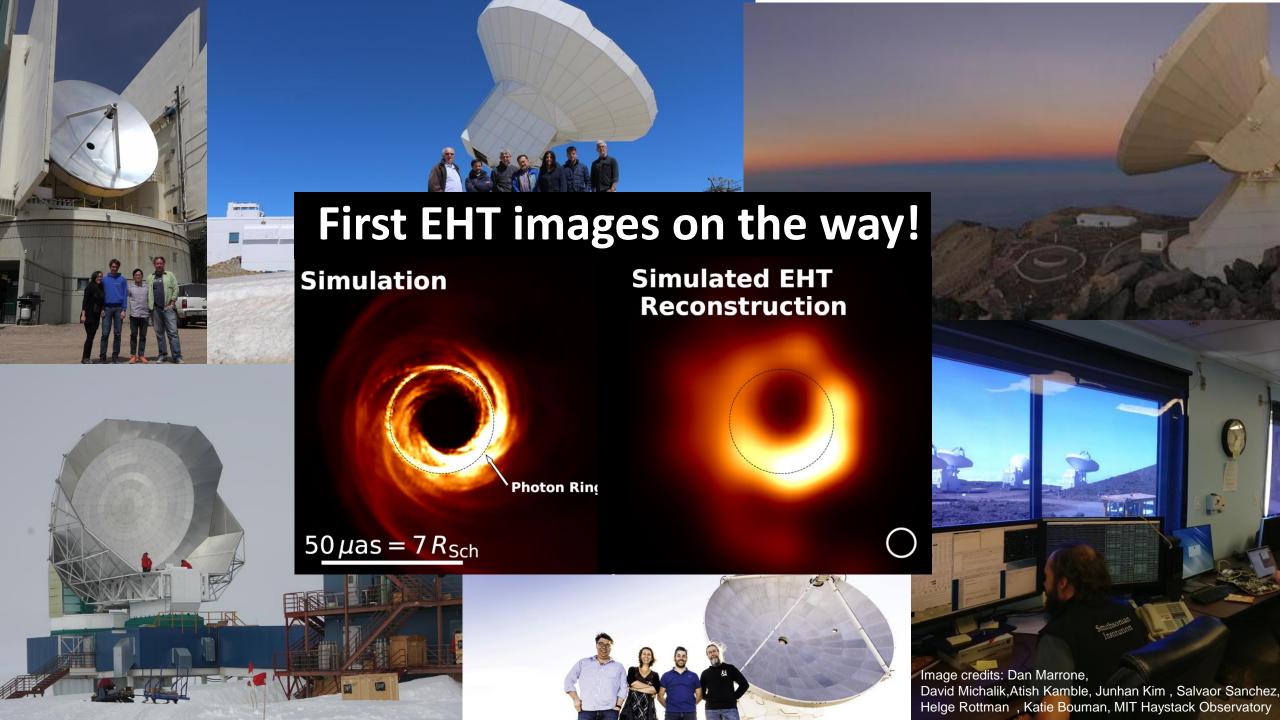
Different plasma heating mechanisms produce qualitatively different images.

• For **Sgr A***:

 Turbulent heating produces a disk-jet structure, which is too anisotropic (when viewed-edge on.)

• For **M87**:

- MAD models produce powerful jets which match VLBI observations.
- But turbulent heating produces too much radiation at the jet base.
- Many features remain unexplained by two-temperature models.
 - Nonthermal electrons.
- EHT images soon!



Thank You!